

Bioengineering 280A
Principles of Biomedical Imaging

Fall Quarter 2004
MRI Lecture 3

TT Liu BE280A, UCSD Fall 2004

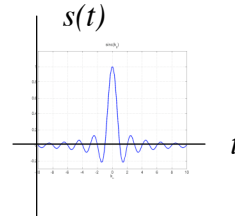
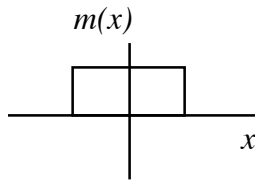
Topics

- Review signal equation
- Sampling requirements
- Slice Selection
- Gradient Echo and Spin Echo
- Image Contrast

TT Liu BE280A, UCSD Fall 2004

MR signal is Fourier Transform

$$\begin{aligned} s(t) &= \int_x \int_y m(x,y) \exp(-j2\pi(k_x(t)x + k_y(t)y)) dx dy \\ &= M(k_x(t), k_y(t)) \\ &= F[m(x,y)]_{k_x(t), k_y(t)} \end{aligned}$$



TT Liu BE280A, UCSD Fall 2004

K-space

At each point in time, the received signal is the Fourier transform of the object

$$s(t) = M(k_x(t), k_y(t)) = F[m(x,y)]_{k_x(t), k_y(t)}$$

evaluated at the spatial frequencies:

$$k_x(t) = \frac{\gamma}{2\pi} \int_0^t G_x(\tau) d\tau$$

$$k_y(t) = \frac{\gamma}{2\pi} \int_0^t G_y(\tau) d\tau$$

Thus, the gradients control our position in k-space. The design of an MRI pulse sequence requires us to efficiently cover enough of k-space to form our image.

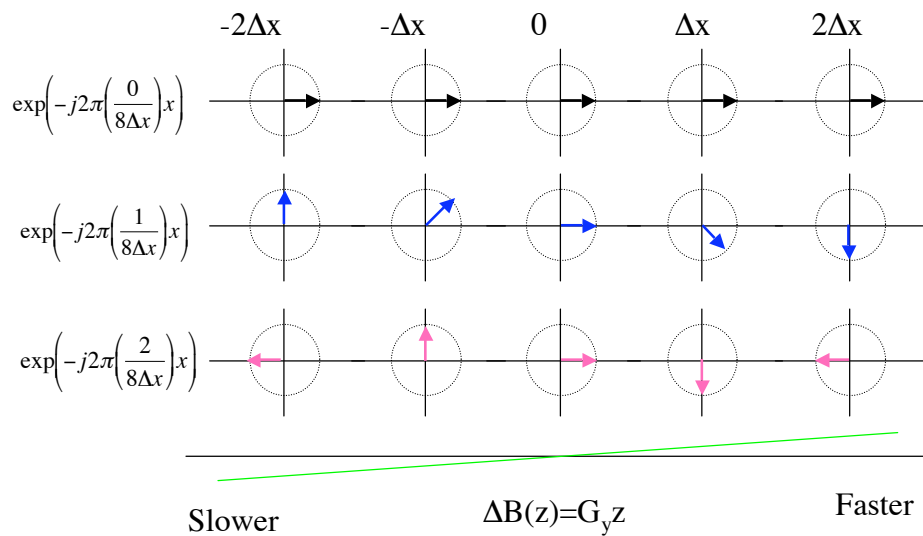
TT Liu BE280A, UCSD Fall 2004

There is nothing that nuclear spins will not do for you, as long as you treat them as human beings.

Erwin Hahn

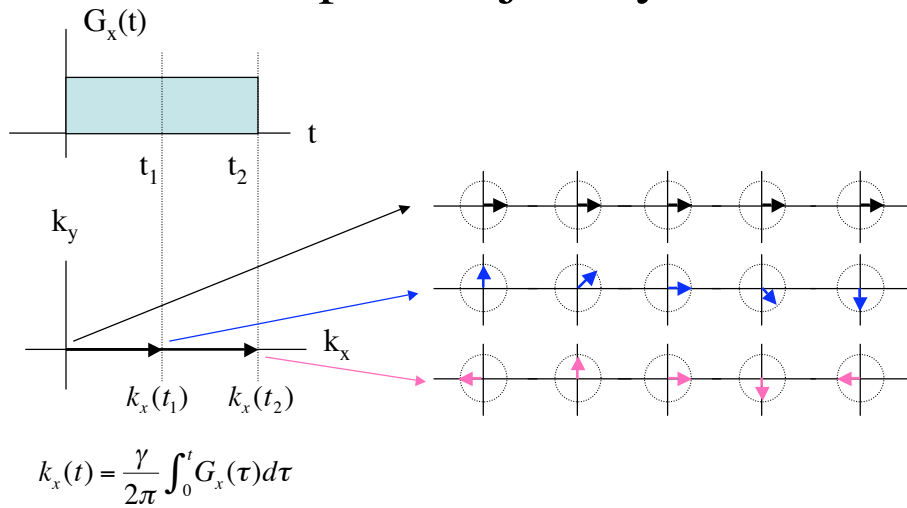
TT Liu BE280A, UCSD Fall 2004

Interpretation



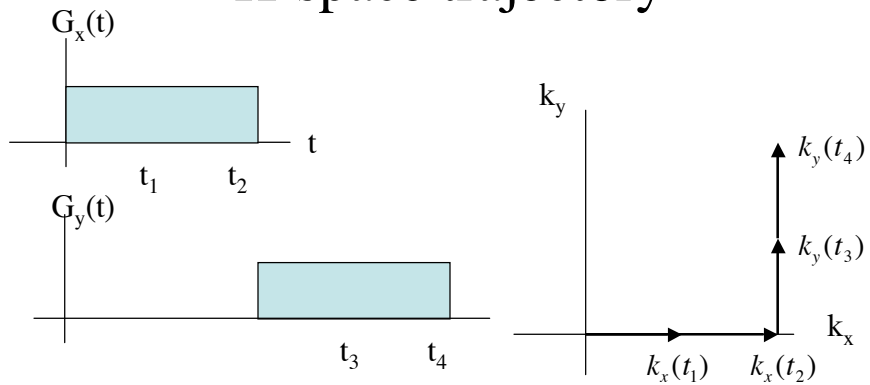
TT Liu BE280A, UCSD Fall 2004

K-space trajectory



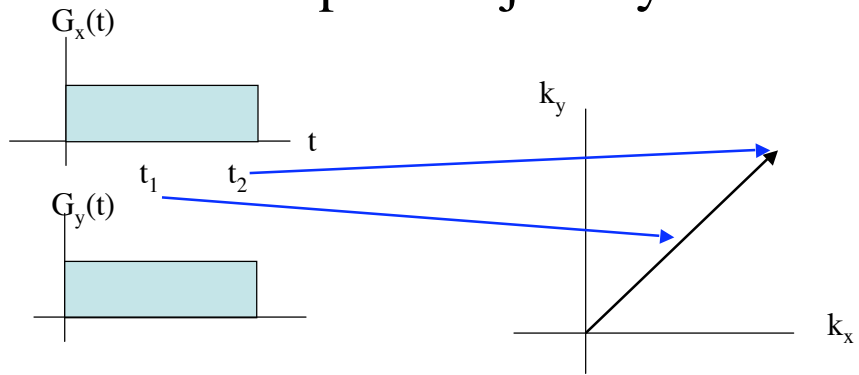
TT Liu BE280A, UCSD Fall 2004

K-space trajectory



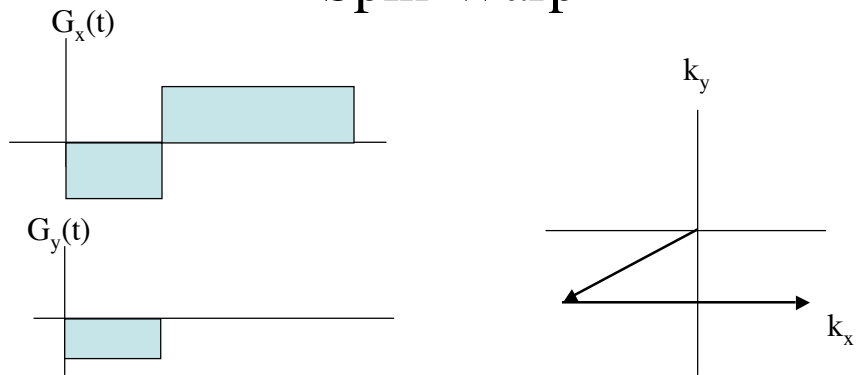
TT Liu BE280A, UCSD Fall 2004

K-space trajectory



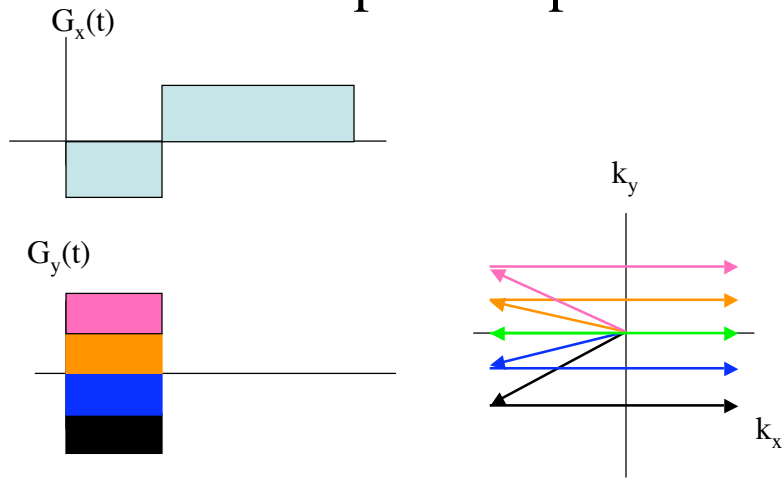
TT Liu BE280A, UCSD Fall 2004

Spin-Warp



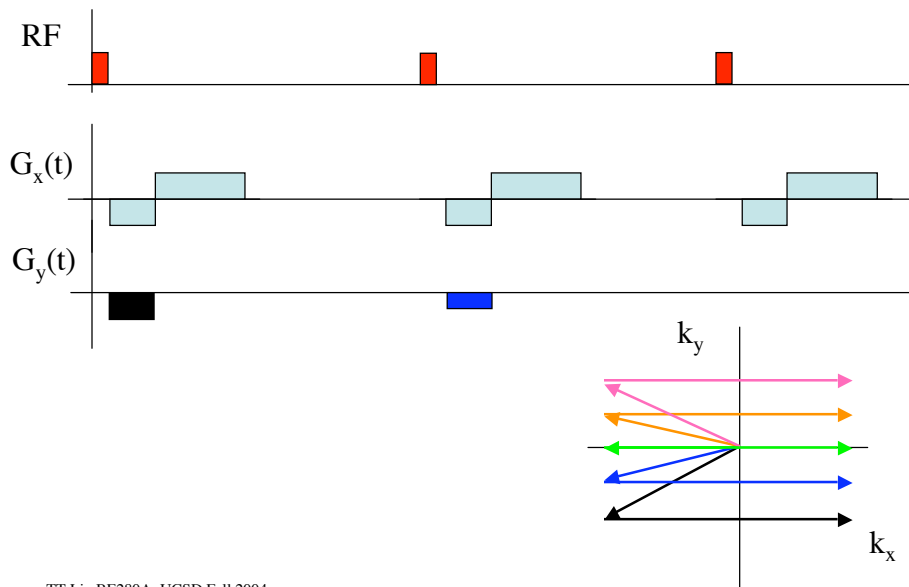
TT Liu BE280A, UCSD Fall 2004

Spin-Warp



TT Liu BE280A, UCSD Fall 2004

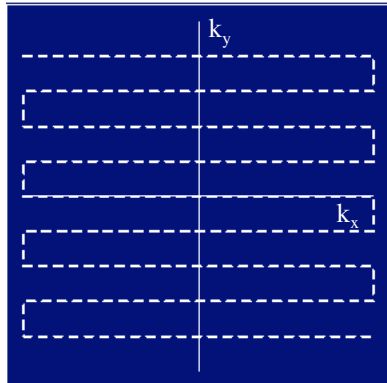
Spin-Warp Pulse Sequence



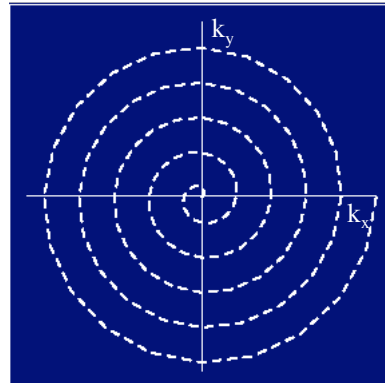
TT Liu BE280A, UCSD Fall 2004

K-space trajectories

EPI



Spiral

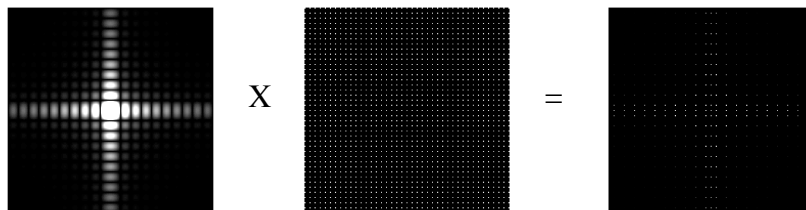


TT Liu BE280A, UCSD Fall 2004

Credit: Larry Frank

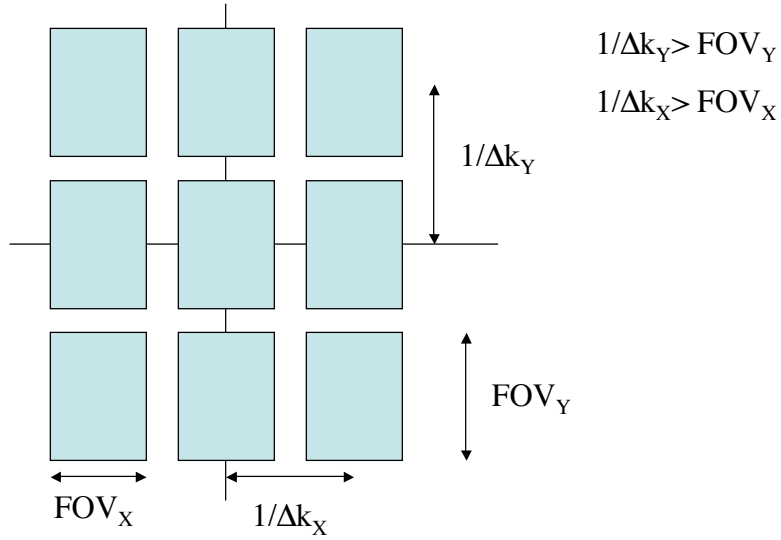


\leftrightarrow
 $1/\Delta k$



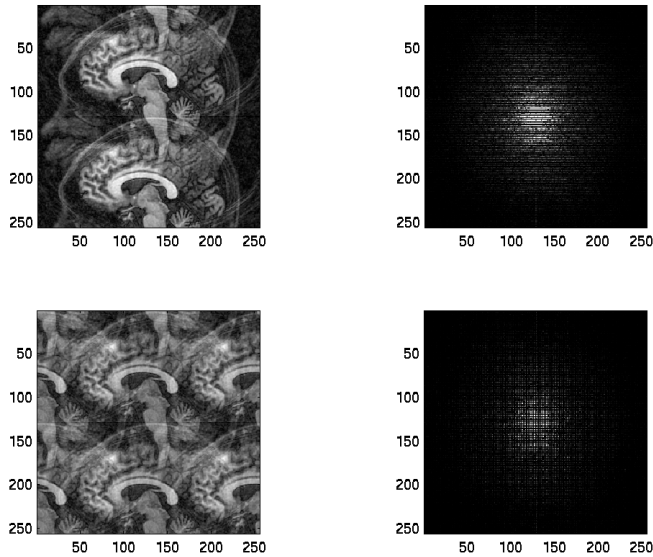
TT Liu BE280A, UCSD Fall 2004

Nyquist Conditions

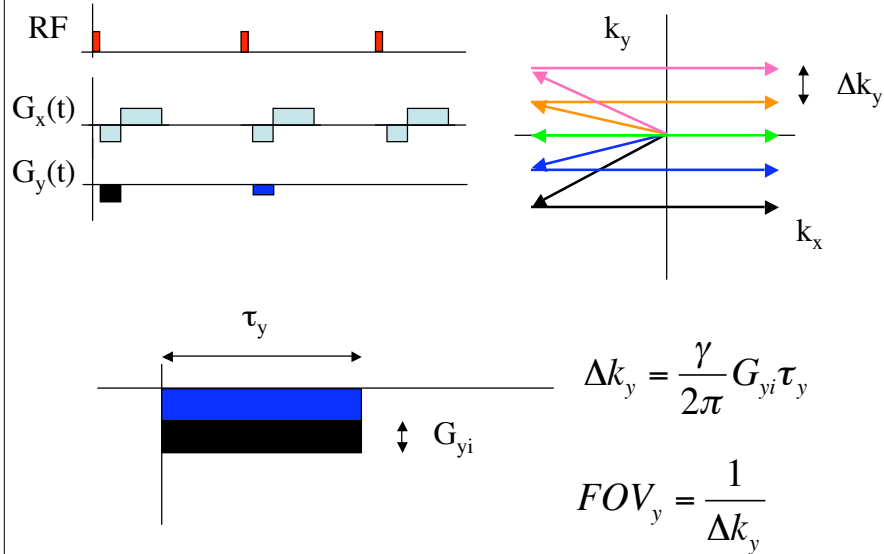


TT Liu BE280A, UCSD Fall 2004

Aliasing

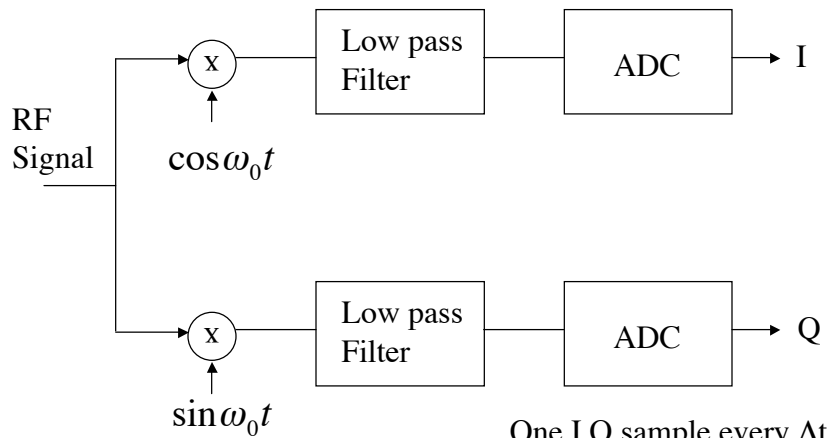


Sampling in k_y



TT Liu BE280A, UCSD Fall 2004

Sampling in k_x

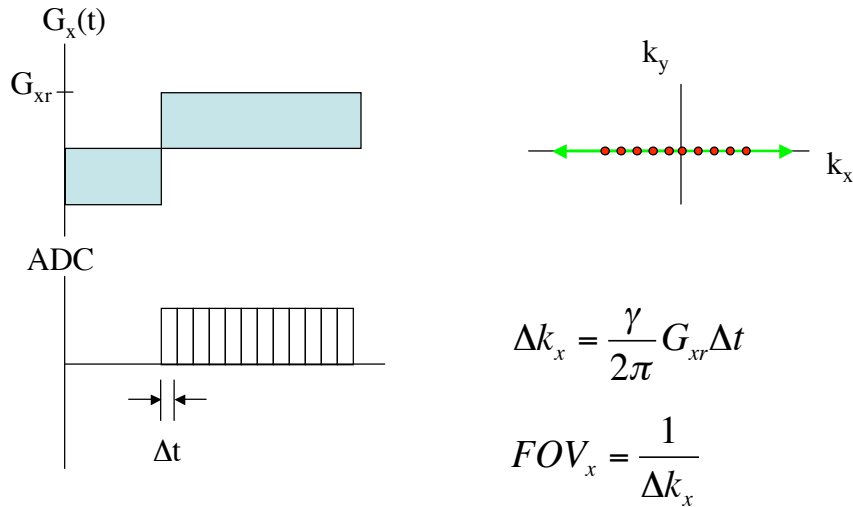


Note: In practice, there are number of ways of implementing this processing.

$$M = I + jQ$$

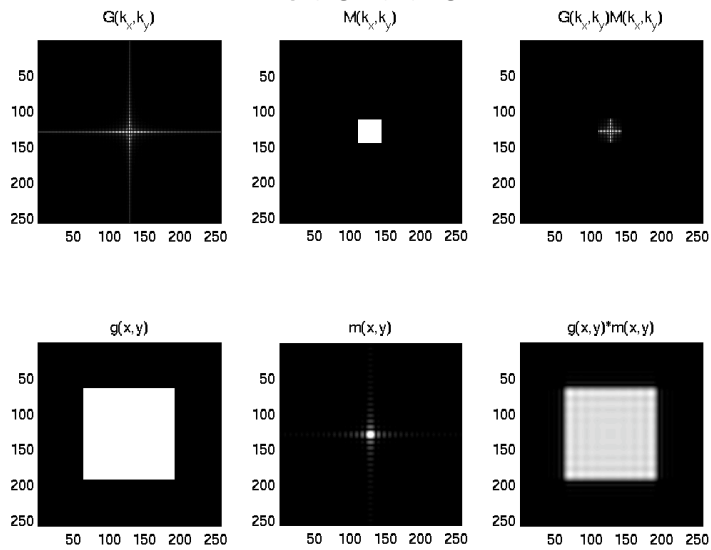
TT Liu BE280A, UCSD Fall 2004

Sampling in k_x



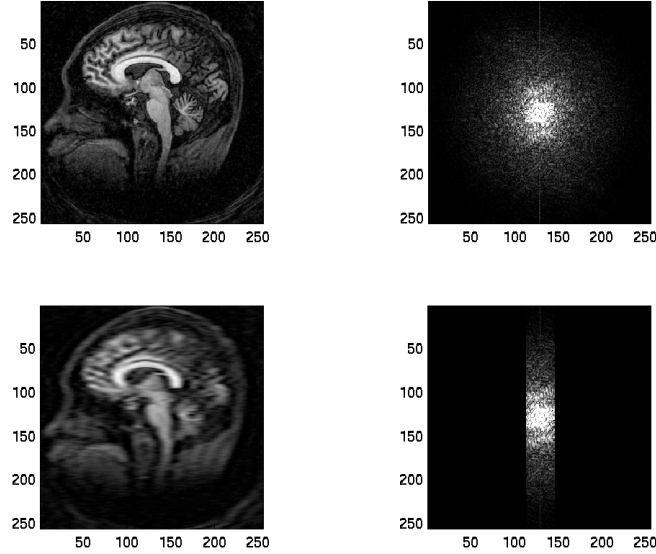
TT Liu BE280A, UCSD Fall 2004

Resolution



TT Liu BE280A, UCSD Fall 2004

Resolution

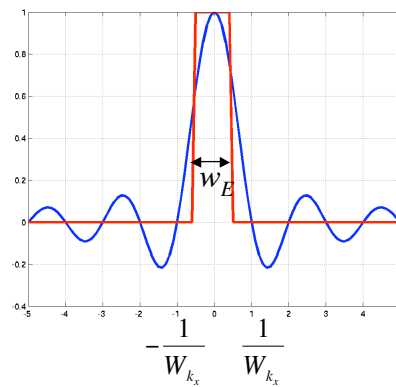


Effective Width

$$w_E = \frac{1}{w(0)} \int_{-\infty}^{\infty} w(x) dx$$

Example

$$\begin{aligned} w_E &= \frac{1}{1} \int_{-\infty}^{\infty} \text{sinc}(W_{k_x} x) dx \\ &= F[\text{sinc}(W_{k_x} x)] \Big|_{k_x=0} \\ &= \frac{1}{W_{k_x}} \text{rect}\left(\frac{k_x}{W_{k_x}}\right) \Big|_{k_x=0} \\ &= \frac{1}{W_{k_x}} \end{aligned}$$

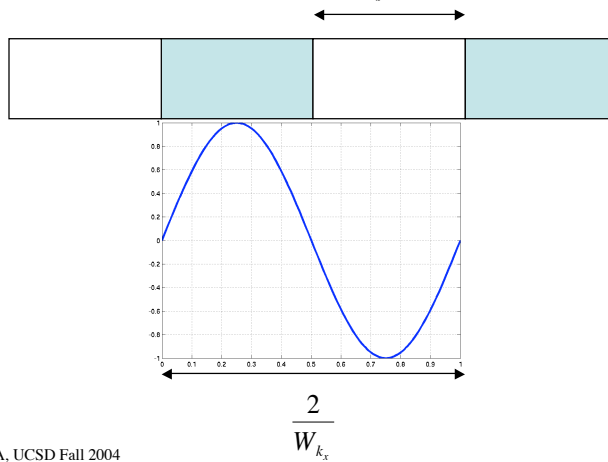


Resolution and spatial frequency

With a window of width W_{k_x} the highest spatial frequency is $W_{k_x}/2$.

This corresponds to a spatial period of $2/W_{k_x}$.

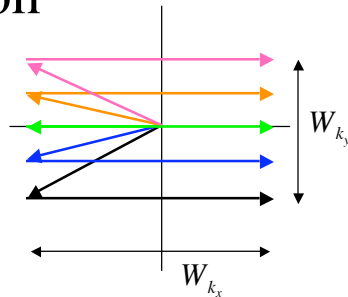
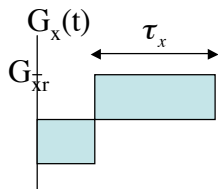
$$\frac{1}{W_{k_x}} = \text{Effective Width} = \delta_x = \text{Resolution}$$



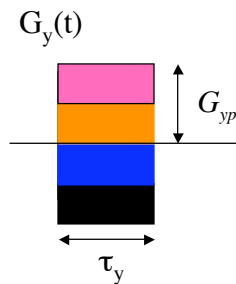
TT Liu BE280A, UCSD Fall 2004

Resolution

$$\delta_x = \frac{1}{W_{k_x}} = \frac{1}{2k_{x,\max}} = \frac{1}{\frac{\gamma}{2\pi} G_{xr} \tau_x}$$



$$\delta_y = \frac{1}{W_{k_y}} = \frac{1}{2k_{y,\max}} = \frac{1}{\frac{\gamma}{2\pi} 2G_{yp} \tau_y}$$



TT Liu BE280A, UCSD Fall 2004

Example

Goal:

$$FOV_x = FOV_y = 25.6 \text{ cm}$$

$$\delta_x = \delta_y = 0.1 \text{ cm}$$

Readout Gradient:

$$FOV_x = \frac{1}{\frac{\gamma}{2\pi} G_{xr} \Delta t}$$

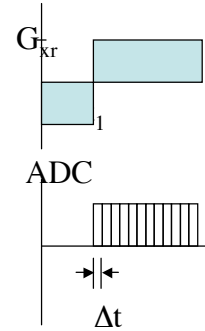
Pick $\Delta t = 32 \text{ } \mu\text{sec}$

$$G_{xr} = \frac{1}{FOV_x \frac{\gamma}{2\pi} \Delta t} = \frac{1}{(25.6 \text{ cm})(42.57 \times 10^6 \text{ T}^{-1} \text{ s}^{-1})(32 \times 10^{-6} \text{ s})}$$

$$= 2.8675 \times 10^{-5} \text{ T/cm}$$

$$= .28675 \text{ G/cm}$$

$$1 \text{ Gauss} = 1 \times 10^{-4} \text{ Tesla}$$



TT Liu BE280A, UCSD Fall 2004

Example

Readout Gradient:

$$\delta_x = \frac{1}{\frac{\gamma}{2\pi} G_{xr} \tau_x}$$

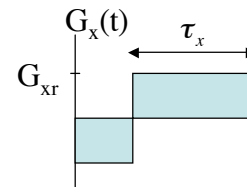
$$\tau_x = \frac{1}{\delta_x \frac{\gamma}{2\pi} G_{xr}} = \frac{1}{(0.1 \text{ cm})(4257 \text{ G}^{-1} \text{ s}^{-1})(0.28675 \text{ G/cm})}$$

$$= 8.192 \text{ ms}$$

$$= N_{\text{read}} \Delta t$$

where

$$N_{\text{read}} = \frac{FOV_x}{\delta_x} = 256$$



TT Liu BE280A, UCSD Fall 2004

Example

Phase - Encode Gradient :

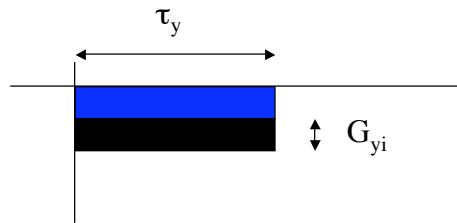
$$FOV_y = \frac{1}{\frac{\gamma}{2\pi} G_{yi} \tau_y}$$

Pick $\tau_y = 4.096 \text{ msec}$

$$G_{yi} = \frac{1}{FOV_y \frac{\gamma}{2\pi} \tau_y} = \frac{1}{(25.6 \text{ cm})(42.57 \times 10^6 \text{ T}^{-1} \text{ s}^{-1})(4.096 \times 10^{-3} \text{ s})}$$

$$= 2.2402 \times 10^{-7} \text{ T/cm}$$

$$= .00224 \text{ G/cm}$$



TT Liu BE280A, UCSD Fall 2004

Example

Phase - Encode Gradient :

$$\delta_y = \frac{1}{\frac{\gamma}{2\pi} 2G_{yp} \tau_y}$$

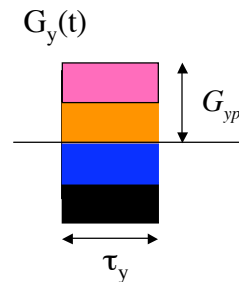
$$G_{yp} = \frac{1}{\delta_y 2 \frac{\gamma}{2\pi} \tau_y} = \frac{1}{(0.1 \text{ cm})(4257 \text{ G}^{-1} \text{ s}^{-1})(4.096 \times 10^{-3} \text{ s})}$$

$$= 0.2868 \text{ G/cm}$$

$$= \frac{N_p}{2} G_{yi}$$

where

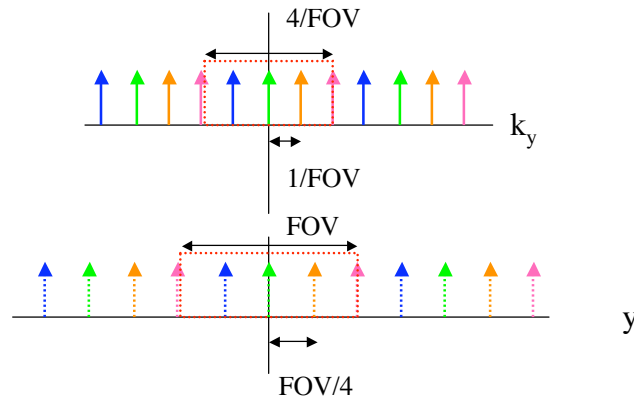
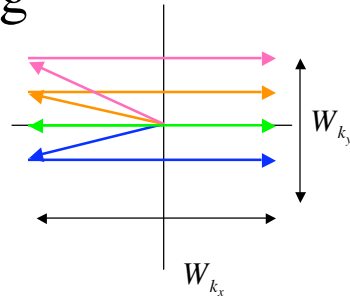
$$N_p = \frac{FOV_y}{\delta_y} = 256$$



TT Liu BE280A, UCSD Fall 2004

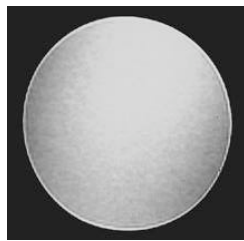
Sampling

In practice, an even number (typically power of 2) sample is usually taken in each direction to take advantage of the Fast Fourier Transform (FFT) for reconstruction.

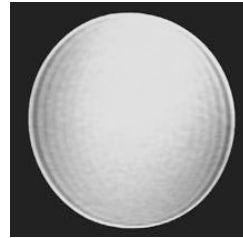


TT Liu BE280A, UCSD Fall 2004

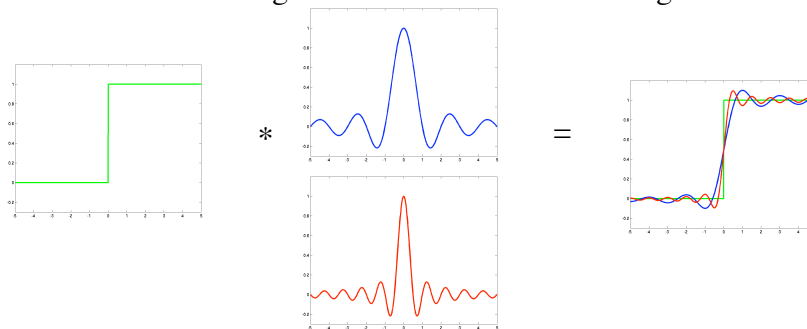
Gibbs Artifact



256x256 image



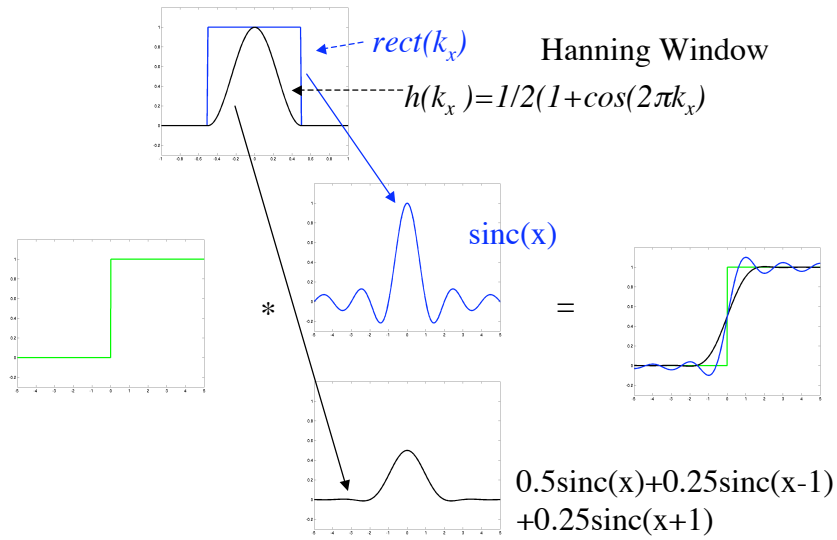
256x128 image



TT Liu BE280A, UCSD Fall 2004

Images from <http://www.mritutor.org/mritutor/gibbs.htm>

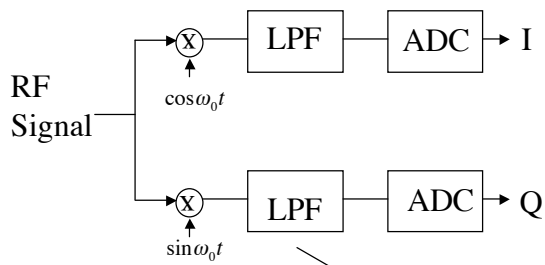
Apodization



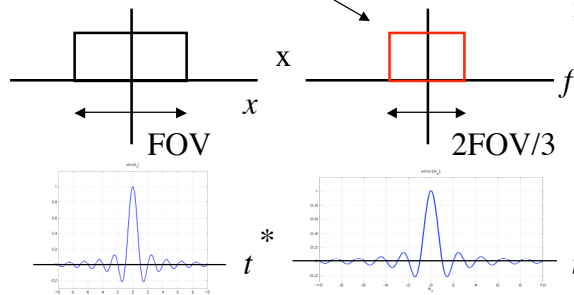
TT Liu BE280A, UCSD Fall 2004

Images from <http://www.mritutor.org/mritutor/gibbs.htm>

Aliasing and Bandwidth

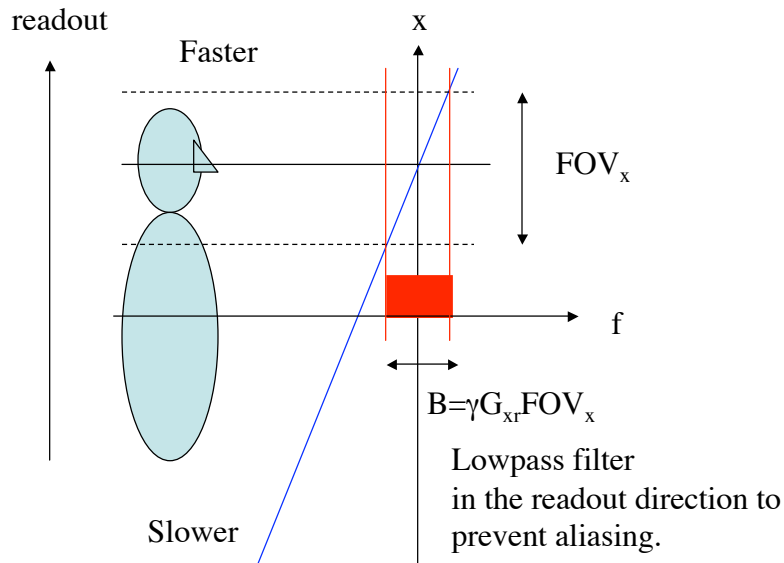


Temporal filtering in the readout direction limits the readout FOV. So there should never be aliasing in the readout direction.



TT Liu BE280A, UCSD Fall 2004

Aliasing and Bandwidth



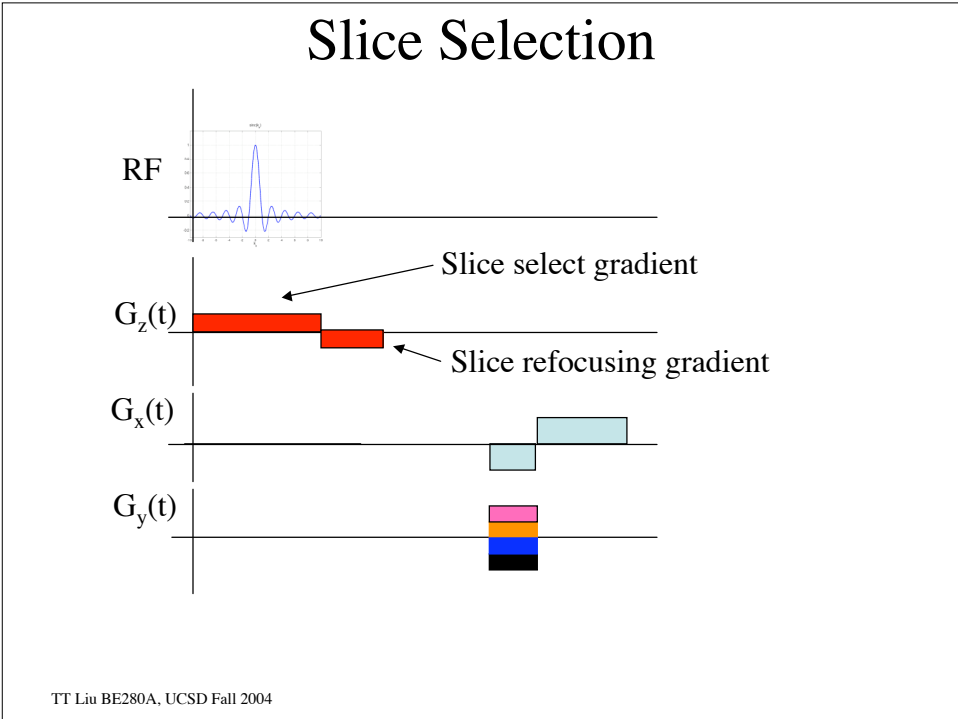
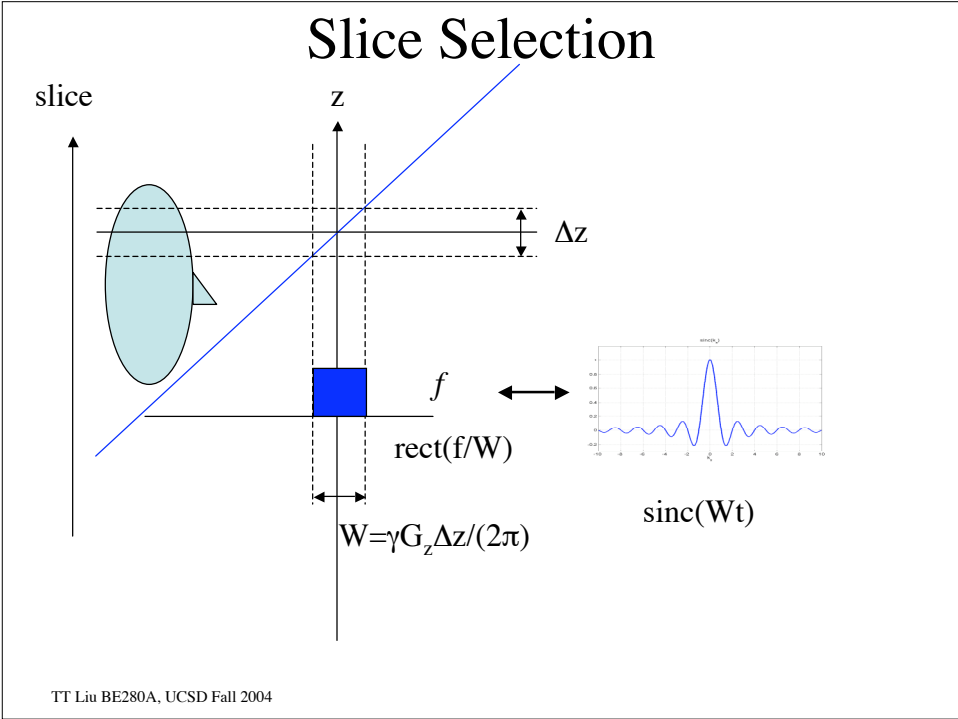
TT Liu BE280A, UCSD Fall 2004

Slice Selection

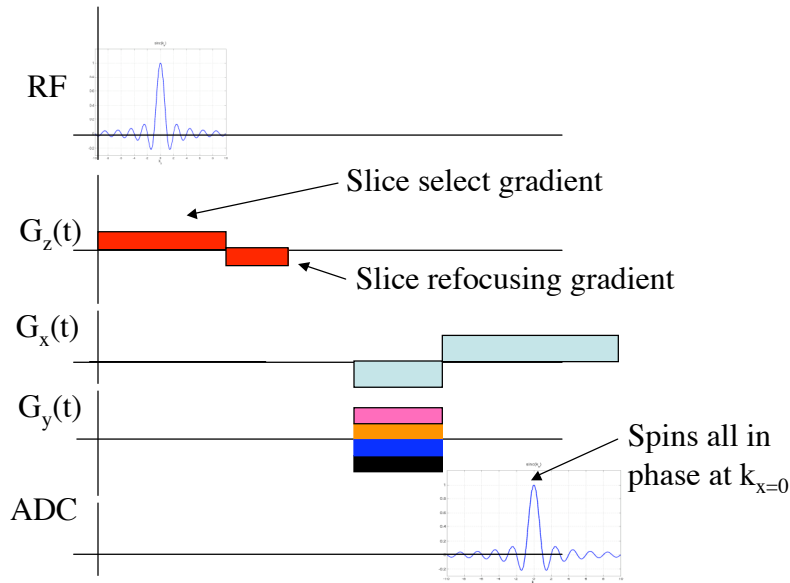
Recall, that we can tip spins away from their equilibrium state by applying a radio-frequency pulse at the Larmor frequency.

In the presence of a spatial gradient G_z , spins in an interval $-\Delta z/2$ to $+\Delta z/2$ have Larmor frequencies ranging from $\omega_0 - \gamma G_z \Delta z/2$ to $\omega_0 + \gamma G_z \Delta z/2$. In order to tip all the spins in this interval, we can apply an RF pulse with energy that is spaced over this frequency interval.

TT Liu BE280A, UCSD Fall 2004



Gradient Echo



TT Liu BE280A, UCSD Fall 2004

Static Inhomogeneities

In the ideal situation, the static magnetic field is totally uniform and the reconstructed object is determined solely by the applied gradient fields. In reality, the magnet is not perfect and will not be totally uniform. Part of this can be addressed by additional coils called "shim" coils, and the process of making the field more uniform is called "shimming". In the old days this was done manually, but modern magnets can do this automatically.

In addition to magnet imperfections, most biological samples are inhomogeneous and this will lead to inhomogeneity in the field. This is because, each tissue has different magnetic properties and will distort the field.

TT Liu BE280A, UCSD Fall 2004

Static Inhomogeneities

The spatial nonuniformity in the field can be modeled by adding an additional term to our signal equation.

$$s_r(t) = \int_V M(\vec{r}, t) dV$$

$$= \int_x \int_y \int_z M(x, y, z, 0) e^{-t/T_2(\vec{r})} e^{-j\omega_0 t} e^{-j\omega_E(\vec{r})t} \exp\left(-j\gamma \int_0^t \vec{G}(\tau) \cdot \vec{r} d\tau\right) dx dy dz$$

The effect of this nonuniformity is to cause the spins to dephase with time and thus for the signal to decrease more rapidly. To first order this can be modeled as an additional decay term of the form

$$s_r(t) = \int_x \int_y \int_z M(x, y, z, 0) e^{-t/T_2(\vec{r})} e^{-t/T_2'(\vec{r})} e^{-j\omega_0 t} \exp\left(-j\gamma \int_0^t \vec{G}(\tau) \cdot \vec{r} d\tau\right) dx dy dz$$

TT Liu BE280A, UCSD Fall 2004

T_2^* decay

The overall decay has the form.

$$\exp\left(-t/T_2^*(\vec{r})\right)$$

where

$$\frac{1}{T_2^*} = \frac{1}{T_2} + \frac{1}{T_2'}$$

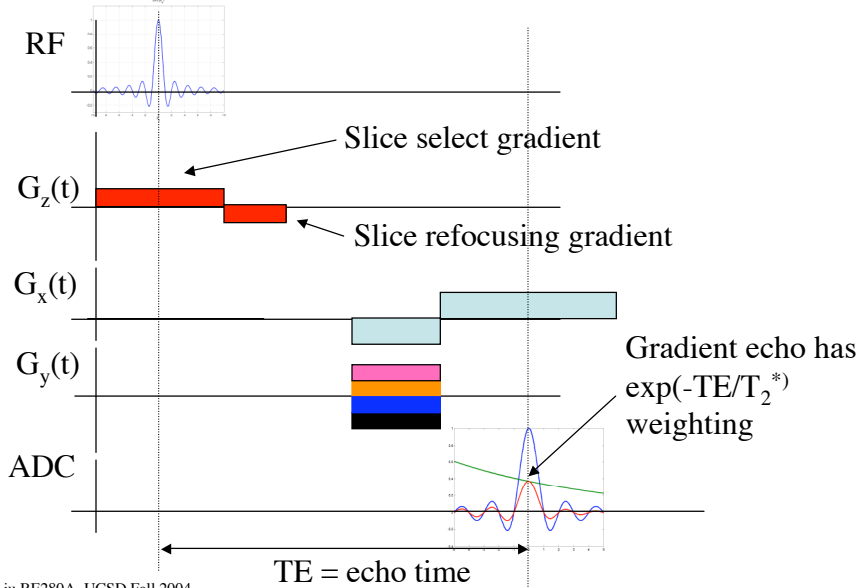
Due to random motions of spins.
Not reversible.

Due to static inhomogeneities. Reversible with a spin-echo sequence.

TT Liu BE280A, UCSD Fall 2004

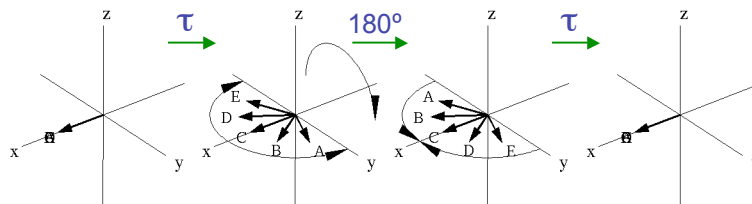
T_2^* decay

Gradient echo sequences exhibit T_2^* decay.



Spin Echo

Discovered by Erwin Hahn in 1950.



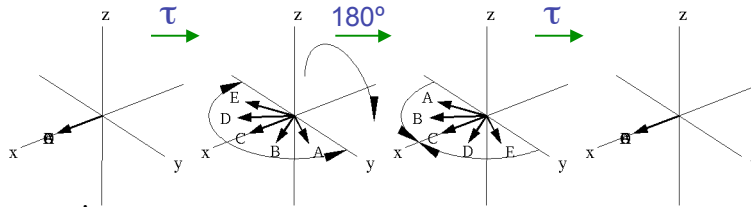
The spin-echo can refocus the dephasing of spins due to static inhomogeneities. However, there will still be T_2 dephasing due to random motion of spins.

There is nothing that nuclear spins will not do for you, as long as you treat them as human beings. Erwin Hahn

TT Liu BE280A, UCSD Fall 2004

Image: Larry Frank

Spin Echo



Phase at time τ

$$\varphi(\tau) = \int_0^\tau -\omega_E(\vec{r}) dt = -\omega_E(\vec{r})\tau$$

Phase after 180 pulse

$$\varphi(\tau^+) = \omega_E(\vec{r})\tau$$

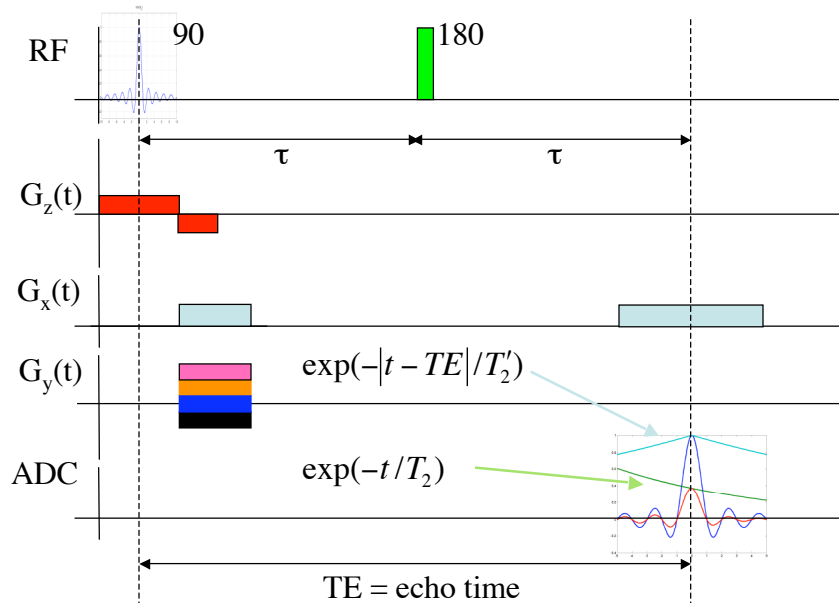
Phase at time 2τ

$$\varphi(2\tau) = -\omega_E(\vec{r})\tau + \omega_E(\vec{r})\tau = 0$$

TT Liu BE280A, UCSD Fall 2004

Image: Larry Frank

Spin Echo Pulse Sequence



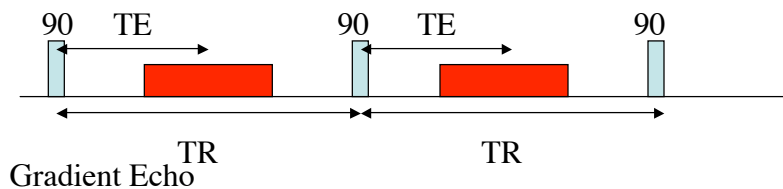
TT Liu BE280A, UCSD Fall 2004

Image Contrast

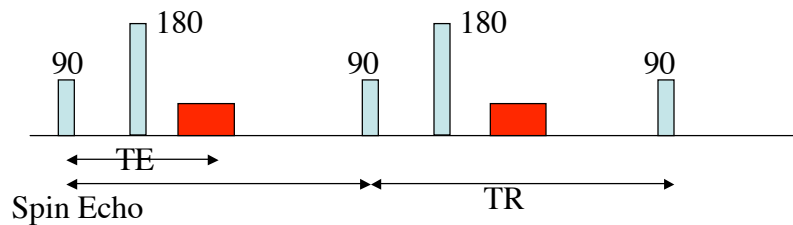
Different tissues exhibit different relaxation rates, T_1 , T_2 , and T_2^* . In addition different tissues can have different densities of protons. By adjusting the pulse sequence, we can create contrast between the tissues. The most basic way of creating contrast is adjusting the two sequence parameters: TE (echo time) and TR (repetition time).

TT Liu BE280A, UCSD Fall 2004

Saturation Recovery Sequence



$$I(x, y) = \rho(x, y) \left[1 - e^{-TR/T_1(x, y)} \right] e^{-TE/T_2^*(x, y)}$$



$$I(x, y) = \rho(x, y) \left[1 - e^{-TR/T_1(x, y)} \right] e^{-TE/T_2(x, y)}$$

TT Liu BE280A, UCSD Fall 2004

T1-Weighted Scans

Make TE very short compared to either T_2 or T_2^* . The resultant image has both proton and T_1 weighting.

$$I(x, y) \approx \rho(x, y) \left[1 - e^{-TR/T_1(x, y)} \right]$$

TT Liu BE280A, UCSD Fall 2004

T2-Weighted Scans

Make TR very long compared to T_1 and use a spin-echo pulse sequence. The resultant image has both proton and T_2 weighting.

$$I(x, y) \approx \rho(x, y) e^{-TE/T_2}$$

TT Liu BE280A, UCSD Fall 2004

Proton Density Weighted Scans

Make TR very long compared to T_1 and use a very short TE. The resultant image is proton density weighted.

$$I(x, y) \approx \rho(x, y)$$

TT Liu BE280A, UCSD Fall 2004

Example



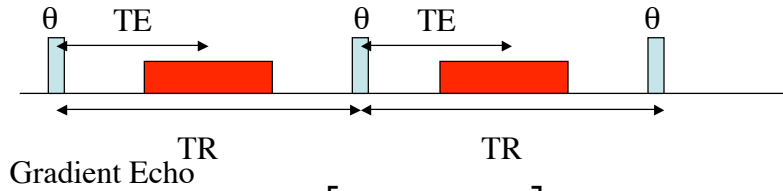
T₁-weighted

Density-weighted

T₂-weighted

TT Liu BE280A, UCSD Fall 2004

FLASH sequence



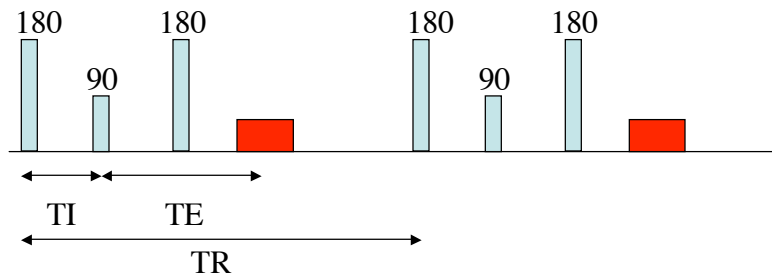
$$I(x, y) = \rho(x, y) \frac{\left[1 - e^{-TR/T_1(x,y)}\right] \sin \theta}{\left[1 - e^{-TR/T_1(x,y)} \cos \theta\right]}$$

Signal intensity is maximized at the Ernst Angle

$$\theta_E = \cos^{-1}(\exp(-TR/T_1))$$

TT Liu BE280A, UCSD Fall 2004

Inversion Recovery



$$I(x, y) = \rho(x, y) \left[1 - 2e^{-TI/T_1(x,y)} + e^{-TR/T_1(x,y)}\right] e^{-TE/T_2(x,y)}$$

Intensity is zero when inversion time is

$$TI = -T_1 \ln \left[\frac{1 + \exp(-TR/T_1)}{2} \right]$$

TT Liu BE280A, UCSD Fall 2004