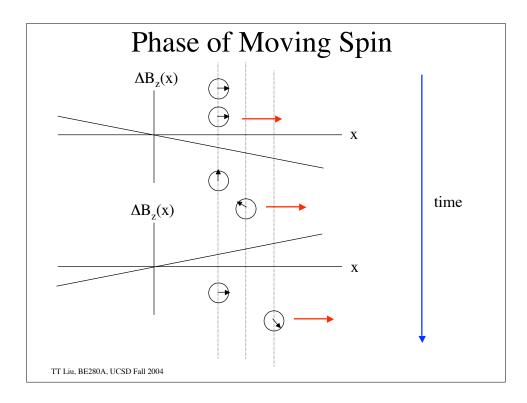


## Moving Spins

So far we have assumed that the spins are not moving (aside from thermal motion giving rise to relaxation), and contrast has been based upon  $T_1$ ,  $T_2$ , and proton density. We were able to achieve different contrasts by adjusting the appropriate pulse sequence parameters.

Biological samples are filled with moving spins, and we can also use MRI to image the movement. Examples: blood flow, diffusion of water in the white matter tracts. In addition, we can also sometimes induce motion into the object to image its mechanical properties, e.g. imaging of stress and strain with MR elastography.

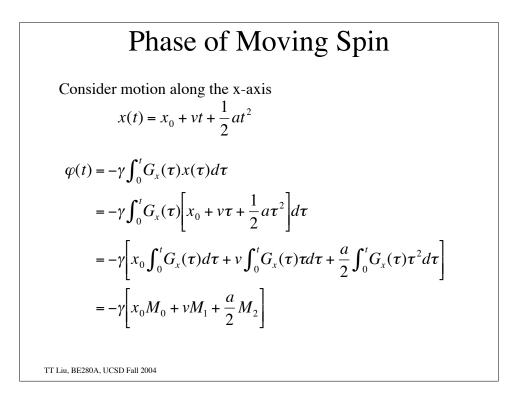
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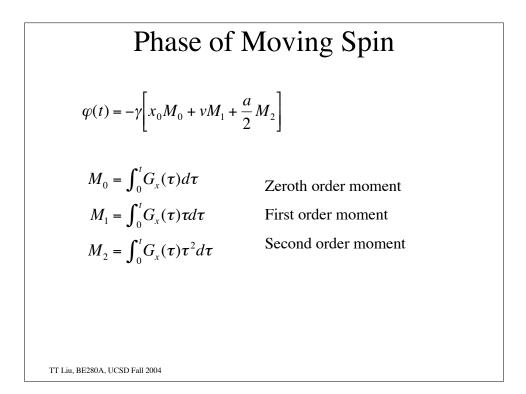


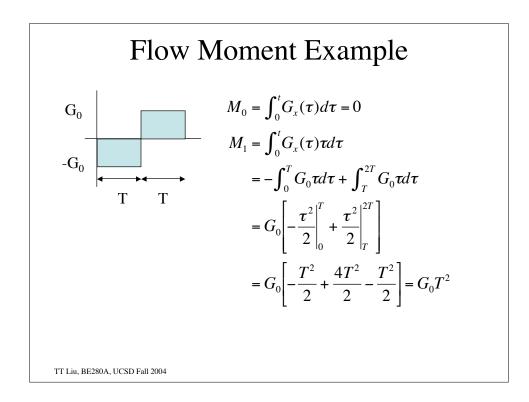
## Phase of a Moving Spin

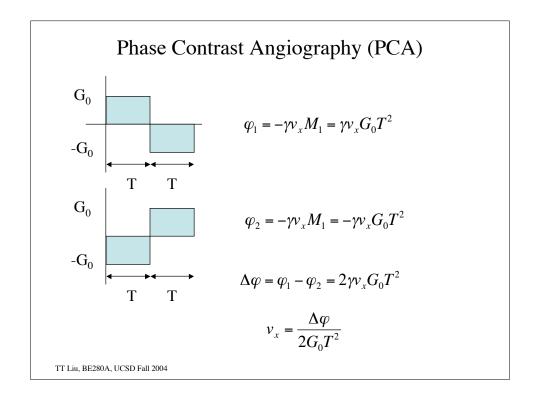
$$\begin{split} \varphi(t) &= -\int_0^t \Delta \omega(\tau) d\tau \\ &= -\int_0^t \gamma \Delta B(\tau) d\tau \\ &= -\int_0^t \gamma \vec{G}(\tau) \cdot \vec{r}(\tau) d\tau \\ &= -\gamma \int_0^t \Big[ G_x(\tau) x(\tau) + G_y(\tau) y(\tau) + G_z(\tau) z(\tau) \Big] d\tau \end{split}$$

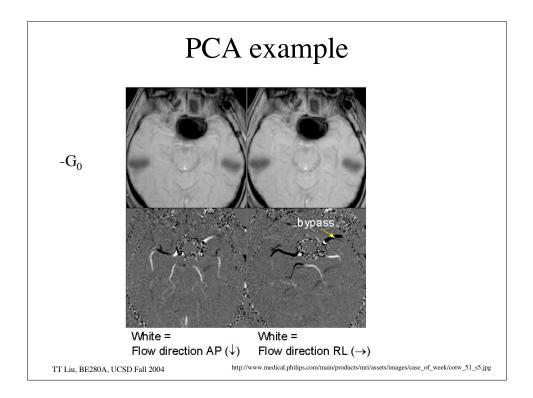
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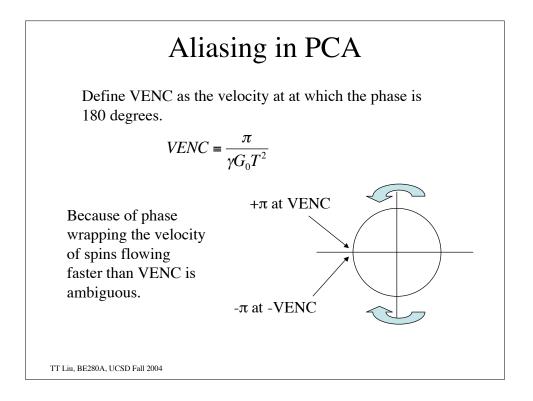


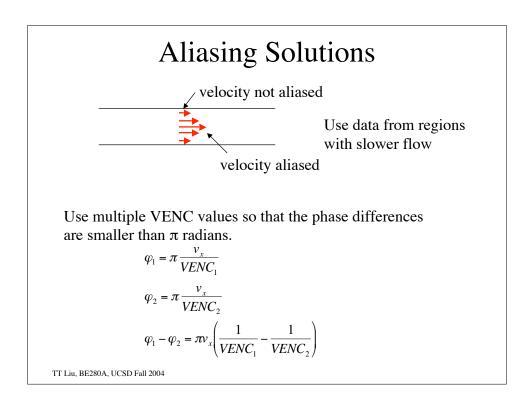












## Velocity k-space

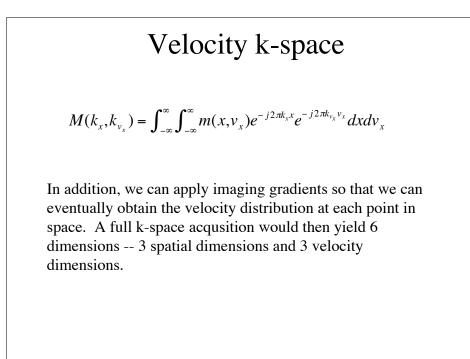
A bipolar gradient introduces a phase modulation across velocities of the form  $\varphi(v_x) = -\gamma v_x G_0 T^2$ 

We can make measurements with different amounts of phase modulation and then integrate over velocities to obtain

$$M(k_{v_x}) = \int_{-\infty}^{\infty} m(v_x) e^{j\varphi(v_x)} dv_x$$
  
= 
$$\int_{-\infty}^{\infty} m(v_x) e^{-j\gamma v_x G_0 T^2} dv_x$$
  
= 
$$\int_{-\infty}^{\infty} m(v_x) e^{-j2\pi k_{v_x} v_x} dv_x$$
  
= 
$$F[m(v_x)] \text{ with } k_{v_x} = \frac{\gamma}{2\pi} G_0 T^2$$

By making measurements with bipolar gradients of varying amplitudes/durations and taking the inverse transform of the measurements, we can obtain the velocity distribution.

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