

Bioengineering 280A  
Principles of Biomedical Imaging

Fall Quarter 2014  
MRI Lecture 2

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## 2D Fourier Transform

Fourier Transform

$$G(k_x, k_y) = F[g(x, y)] = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} g(x, y) e^{-j2\pi(k_x x + k_y y)} dx dy$$

Inverse Fourier Transform

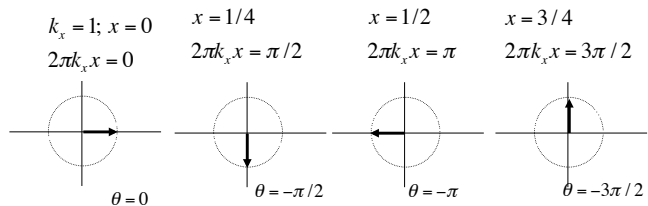
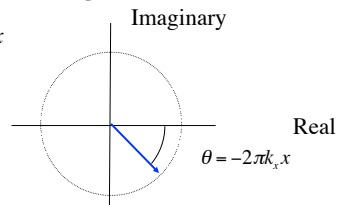
$$g(x, y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} G(k_x, k_y) e^{j2\pi(k_x x + k_y y)} dk_x dk_y$$

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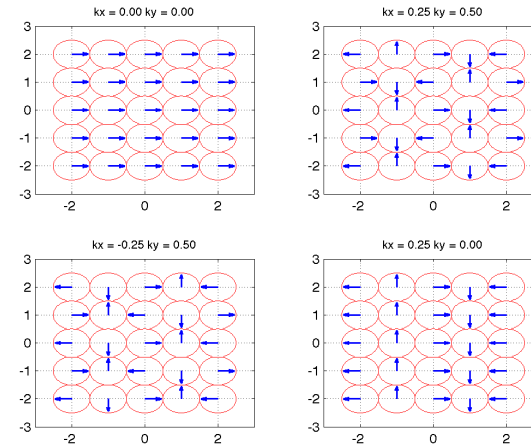
## Phasor Diagram

$$G(k_x) = \int_{-\infty}^{\infty} g(x) \exp(-j2\pi k_x x) dx$$

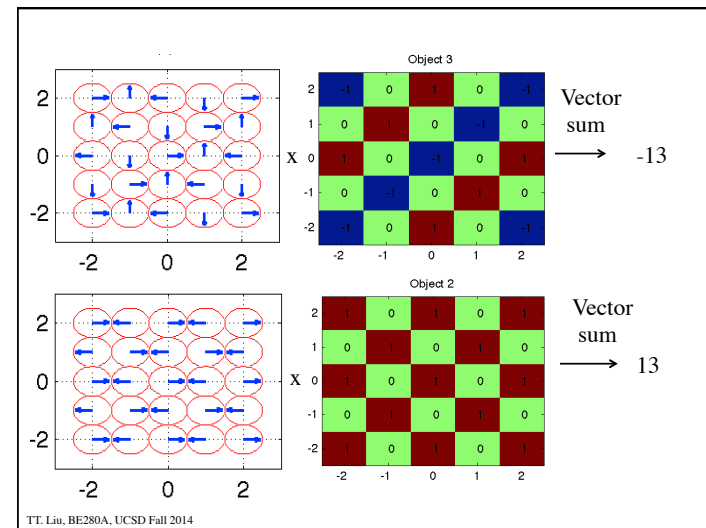
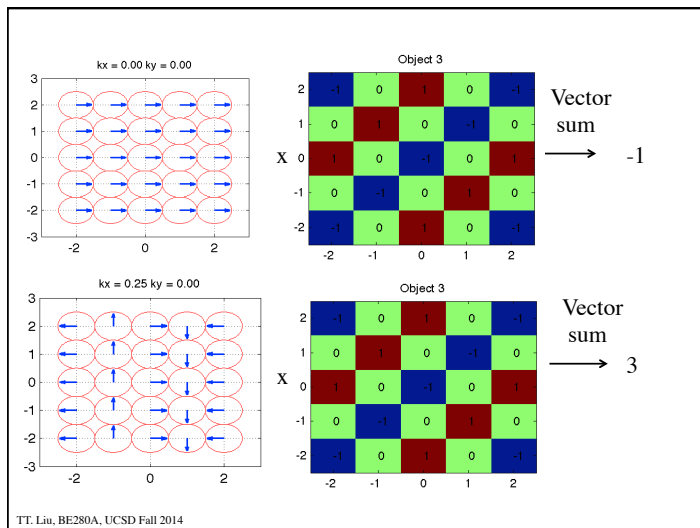
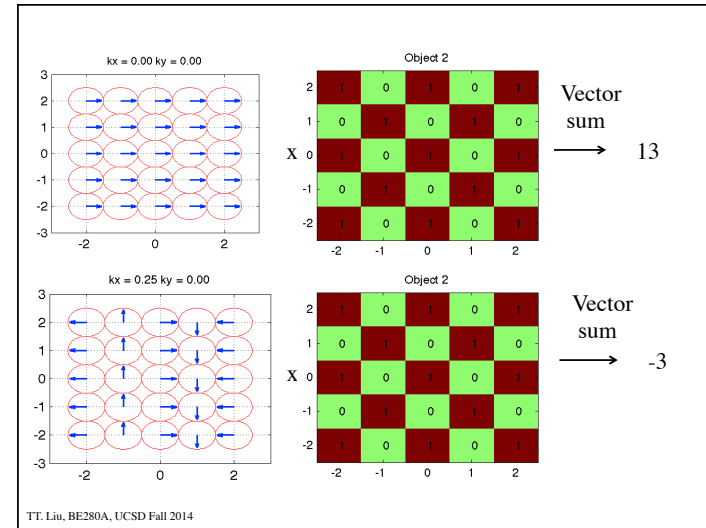
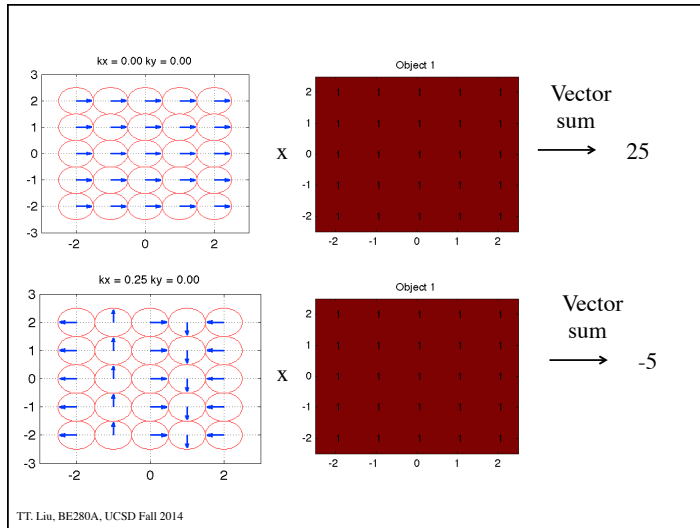
$\theta = -2\pi k_x x$



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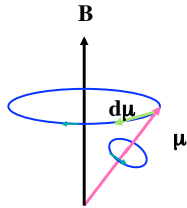


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## Precession

$$\frac{d\boldsymbol{\mu}}{dt} = \boldsymbol{\mu} \times \gamma \mathbf{B}$$



Analogous to motion of a gyroscope

Precesses at an angular frequency of

$$\omega = \gamma B$$

This is known as the **Larmor** frequency.



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## Larmor Frequency

$\omega = \gamma B$  Angular frequency in rad/sec

$f = \gamma B / (2\pi)$  Frequency in cycles/sec or Hertz, Abbreviated Hz

For a 1.5 T system, the Larmor frequency is 63.86 MHz which is 63.86 million cycles per second. For comparison, KPBS-FM transmits at 89.5 MHz.

Note that the earth's magnetic field is about 50  $\mu$ T, so that a 1.5T system is about 30,000 times stronger.

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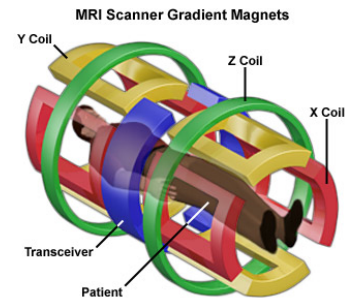
## Gradients

Spins precess at the Larmor frequency, which is proportional to the local magnetic field. In a constant magnetic field  $B_z = B_0$ , all the spins precess at the same frequency (ignoring chemical shift).

Gradient coils are used to add a spatial variation to  $B_z$  such that  $B_z(x,y,z) = B_0 + \Delta B_z(x,y,z)$ . Thus, spins at different physical locations will precess at different frequencies.

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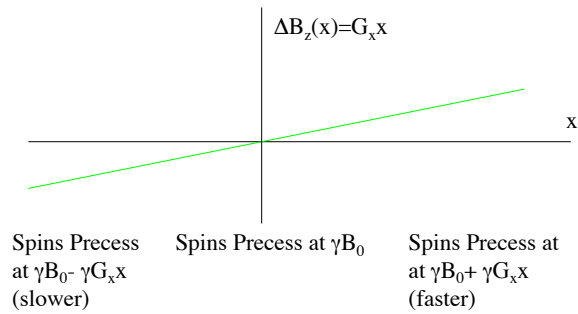
## MRI Gradients



<http://www.magnet.fsu.edu/education/tutorials/magnetacademy/mri/fullarticle.html>  
<http://gulfnnews.com/business/economy/german-factories-race-to-meet-big-surge-in-orders-1.638066>

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## Interpretation



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## Rotating Frame of Reference

Reference everything to the magnetic field at isocenter.



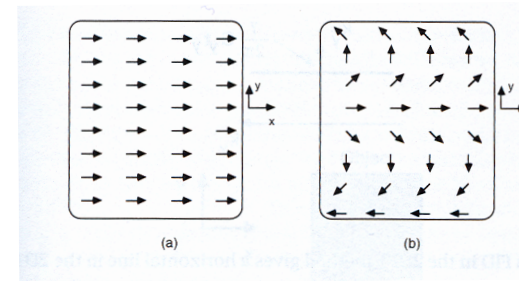
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## Spins



*There is nothing that nuclear spins will not do for you, as long as you treat them as human beings.*  
Erwin Hahn

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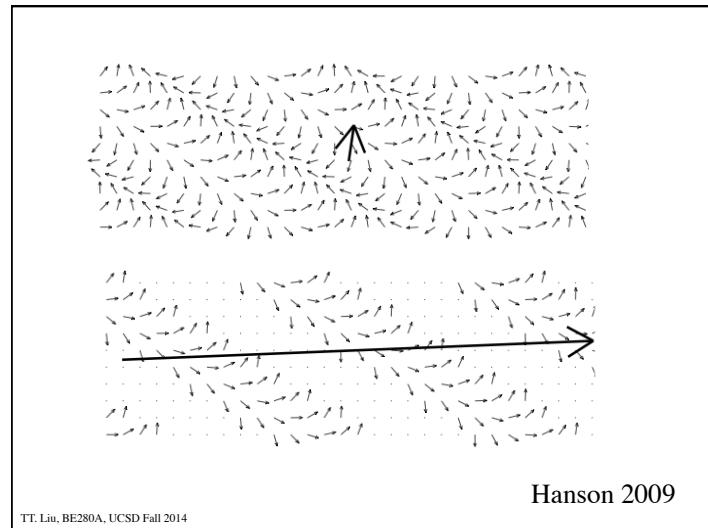
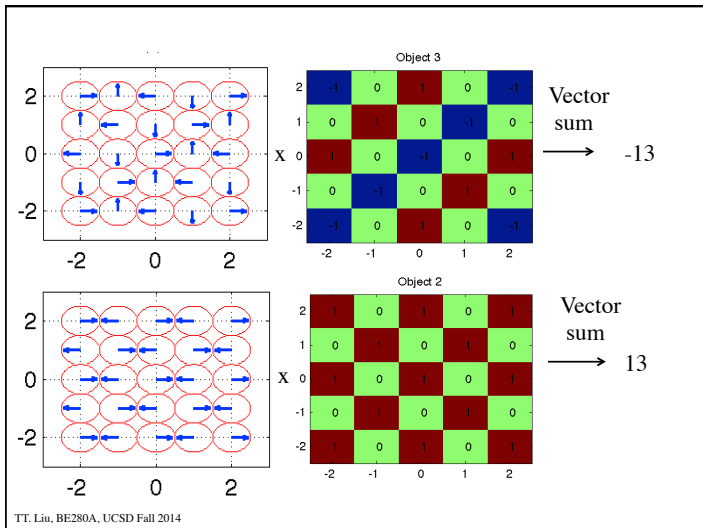
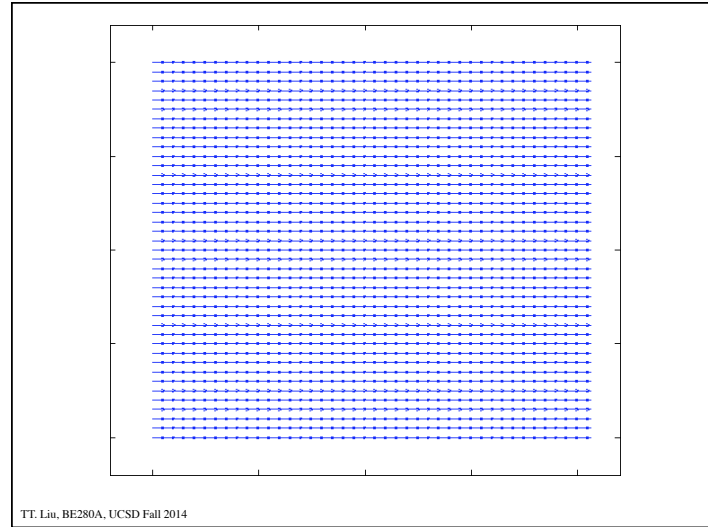
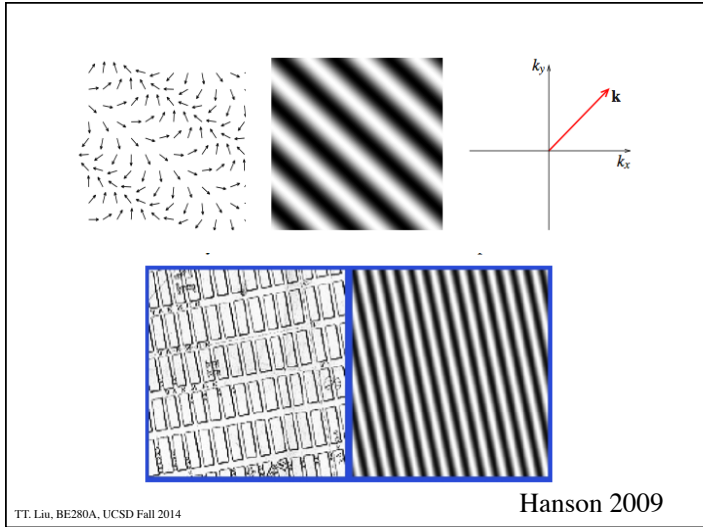


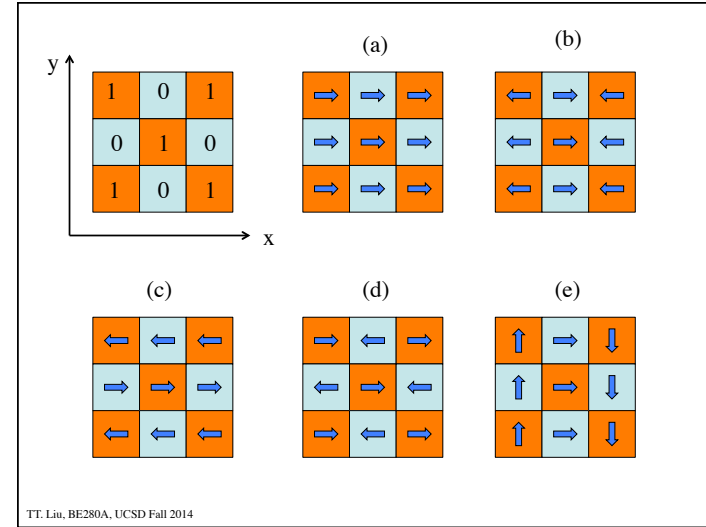
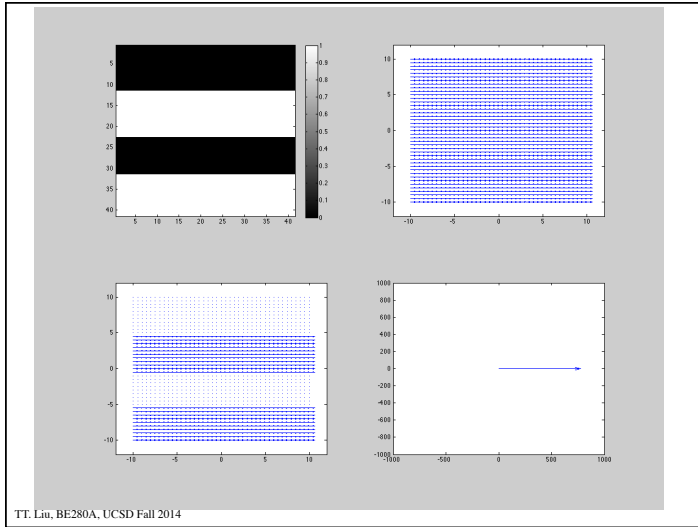
$k_x=0; k_y=0$

$k_x=0; k_y \neq 0$

Fig 3.12 from Nishimura

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### Gradient Fields

$$B_z(x, y, z) = B_0 + \frac{\partial B_z}{\partial x} x + \frac{\partial B_z}{\partial y} y + \frac{\partial B_z}{\partial z} z$$

$$= B_0 + G_x x + G_y y + G_z z$$

$$G_z = \frac{\partial B_z}{\partial z} > 0$$

$$G_y = \frac{\partial B_z}{\partial y} > 0$$

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### Gradient Fields

Define

$$\vec{G} \equiv G_x \hat{i} + G_y \hat{j} + G_z \hat{k} \quad \vec{r} \equiv x \hat{i} + y \hat{j} + z \hat{k}$$

So that

$$G_x x + G_y y + G_z z = \vec{G} \cdot \vec{r}$$

Also, let the gradient fields be a function of time. Then the z-directed magnetic field at each point in the volume is given by :

$$B_z(\vec{r}, t) = B_0 + \vec{G}(t) \cdot \vec{r}$$

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## Time-Varying Gradient Fields

In the presence of time-varying gradients the frequency as a function of space and time is:

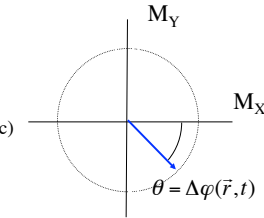
$$\begin{aligned}\omega(\vec{r}, t) &= \gamma B_z(\vec{r}, t) \\ &= \gamma B_0 + \gamma \vec{G}(t) \cdot \vec{r} \\ &= \omega_0 + \Delta\omega(\vec{r}, t)\end{aligned}$$

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## Phase

$$M(\vec{r}, t) = M(\vec{r}, 0)e^{i\varphi(\vec{r}, t)}$$

Phase = angle of the magnetization phasor  
Frequency = rate of change of angle (e.g. radians/sec)  
Phase = time integral of frequency



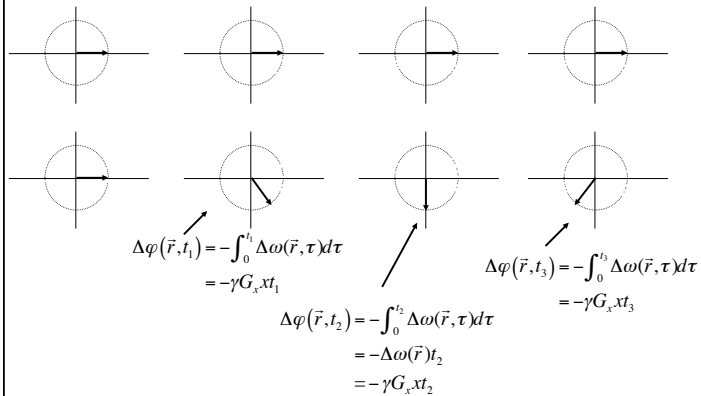
$$\begin{aligned}\varphi(\vec{r}, t) &= -\int_0^t \omega(\vec{r}, \tau) d\tau \\ &= -\omega_0 t + \Delta\varphi(\vec{r}, t)\end{aligned}$$

Where the incremental phase due to the gradients is

$$\begin{aligned}\Delta\varphi(\vec{r}, t) &= -\int_0^t \Delta\omega(\vec{r}, \tau) d\tau \\ &= -\int_0^t \gamma \vec{G}(\vec{r}, \tau) \cdot \vec{r} d\tau\end{aligned}$$

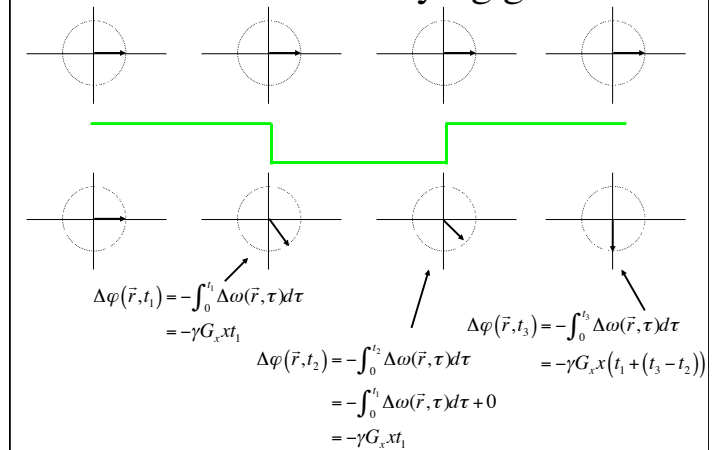
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## Phase with constant gradient



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## Phase with time-varying gradient



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## Time-Varying Gradient Fields

The transverse magnetization is then given by

$$\begin{aligned} M(\vec{r}, t) &= M(\vec{r}, 0) e^{-t/T_2(\vec{r})} e^{i\phi(\vec{r}, t)} \\ &= M(\vec{r}, 0) e^{-t/T_2(\vec{r})} e^{-j\omega_0 t} \exp\left(-j \int_0^t \Delta\omega(\vec{r}, \tau) d\tau\right) \\ &= M(\vec{r}, 0) e^{-t/T_2(\vec{r})} e^{-j\omega_0 t} \exp\left(-j\gamma \int_0^t \vec{G}(\tau) \cdot \vec{r} d\tau\right) \end{aligned}$$

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## Signal Equation

Signal from a volume

$$\begin{aligned} s_r(t) &= \int_V M(\vec{r}, t) dV \\ &= \int_x \int_y \int_z M(x, y, z, 0) e^{-t/T_2(\vec{r})} e^{-j\omega_0 t} \exp\left(-j\gamma \int_0^t \vec{G}(\tau) \cdot \vec{r} d\tau\right) dx dy dz \end{aligned}$$

For now, consider signal from a slice along  $z$  and drop the  $T_2$  term. Define  $m(x, y) = \int_{z_0-\Delta z/2}^{z_0+\Delta z/2} M(\vec{r}, t) dz$

To obtain

$$s_r(t) = \int_x \int_y m(x, y) e^{-j\omega_0 t} \exp\left(-j\gamma \int_0^t \vec{G}(\tau) \cdot \vec{r} d\tau\right) dx dy$$

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## Signal Equation

Demodulate the signal to obtain

$$\begin{aligned} s(t) &= e^{j\omega_0 t} s_r(t) \\ &= \int_x \int_y m(x, y) \exp\left(-j\gamma \int_0^t \vec{G}(\tau) \cdot \vec{r} d\tau\right) dx dy \\ &= \int_x \int_y m(x, y) \exp\left(-j\gamma \int_0^t [G_x(\tau)x + G_y(\tau)y] d\tau\right) dx dy \\ &= \int_x \int_y m(x, y) \exp\left(-j2\pi(k_x(t)x + k_y(t)y)\right) dx dy \end{aligned}$$

Where

$$\begin{aligned} k_x(t) &= \frac{\gamma}{2\pi} \int_0^t G_x(\tau) d\tau \\ k_y(t) &= \frac{\gamma}{2\pi} \int_0^t G_y(\tau) d\tau \end{aligned}$$

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## MR signal is Fourier Transform

$$\begin{aligned} s(t) &= \int_x \int_y m(x, y) \exp\left(-j2\pi(k_x(t)x + k_y(t)y)\right) dx dy \\ &= M(k_x(t), k_y(t)) \\ &= F[m(x, y)] \Big|_{k_x(t), k_y(t)} \end{aligned}$$

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## Recap

- Frequency = rate of change of phase.
- Higher magnetic field -> higher Larmor frequency -> phase changes more rapidly with time.
- With a constant gradient  $G_x$ , spins at different x locations precess at different frequencies -> spins at greater x-values change phase more rapidly.
- With a constant gradient, distribution of phases across x locations changes with time. (phase modulation)
- More rapid change of phase with x -> higher spatial frequency  $k_x$

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## K-space

At each point in time, the received signal is the Fourier transform of the object

$$s(t) = M(k_x(t), k_y(t)) = F[m(x, y)]_{k_x(t), k_y(t)}$$

evaluated at the spatial frequencies:

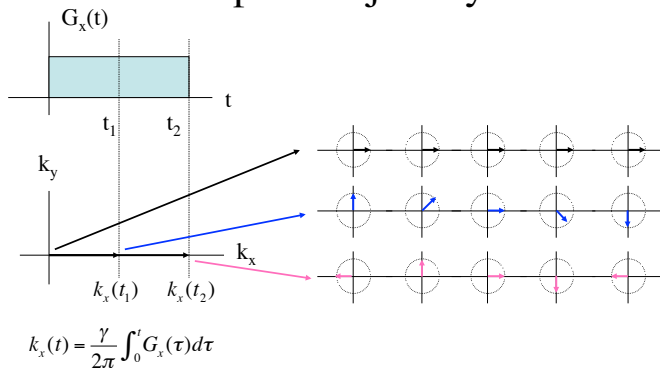
$$k_x(t) = \frac{\gamma}{2\pi} \int_0^t G_x(\tau) d\tau$$

$$k_y(t) = \frac{\gamma}{2\pi} \int_0^t G_y(\tau) d\tau$$

Thus, the gradients control our position in k-space. The design of an MRI pulse sequence requires us to efficiently cover enough of k-space to form our image.

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## K-space trajectory



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## Units

Spatial frequencies ( $k_x, k_y$ ) have units of 1/distance. Most commonly, 1/cm

Gradient strengths have units of (magnetic field)/distance. Most commonly G/cm or mT/m

$\gamma/(2\pi)$  has units of Hz/G or Hz/Tesla.

$$\begin{aligned} k_x(t) &= \frac{\gamma}{2\pi} \int_0^t G_x(\tau) d\tau \\ &= [\text{Hz} / \text{Gauss}] [\text{Gauss} / \text{cm}] [\text{sec}] \\ &= [1 / \text{cm}] \end{aligned}$$

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