

ECE187
Introduction to Biomedical Imaging

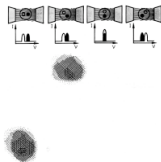
Fall Quarter 2004
MRI Lecture

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History of MRI



1946: Felix Bloch (Stanford) and Edward Purcell (Harvard) demonstrate nuclear magnetic resonance (NMR)



1973: Paul Lauterbur (SUNY) published first MRI image in Nature.

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History of MRI

Late 1970's: First human MRI images

Early 1980's: First commercial MRI systems

1993: functional MRI in humans demonstrated

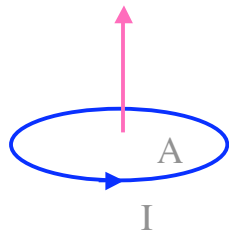
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Spin

- Intrinsic angular momentum of elementary particles -- electrons, protons, neutrons.
- Spin is quantized. Key concept in Quantum Mechanics.

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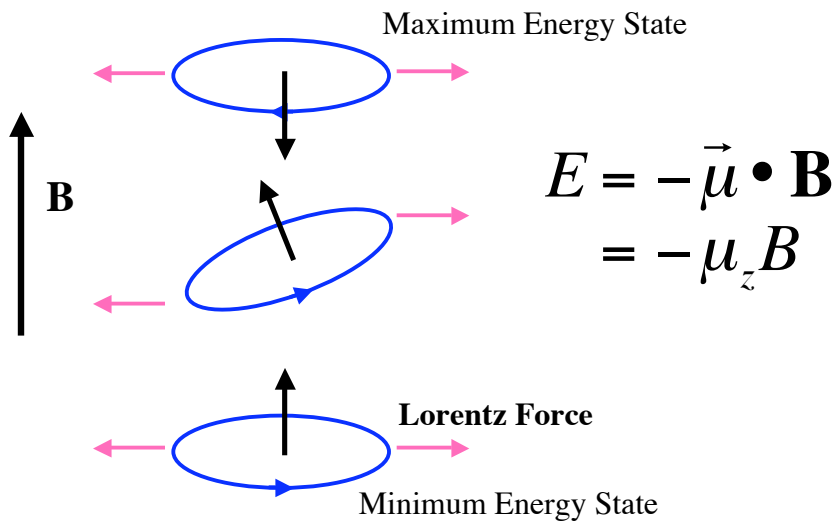
Classical Magnetic Moment



$$\vec{\mu} = IA\hat{n}$$

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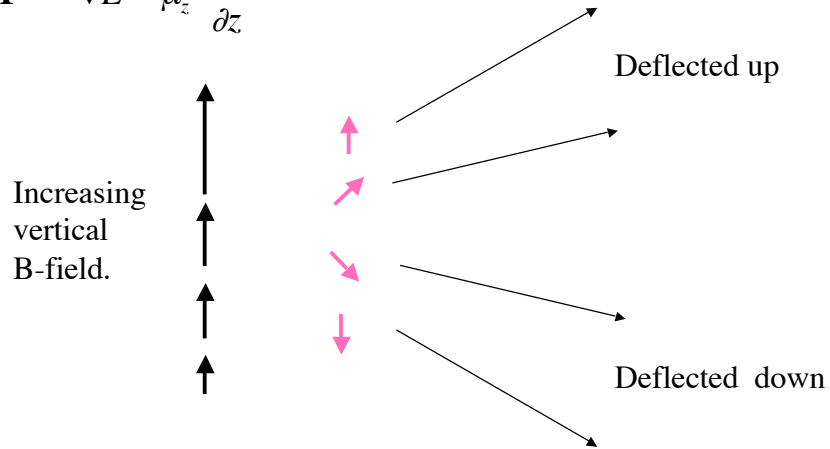
Energy in a Magnetic Field



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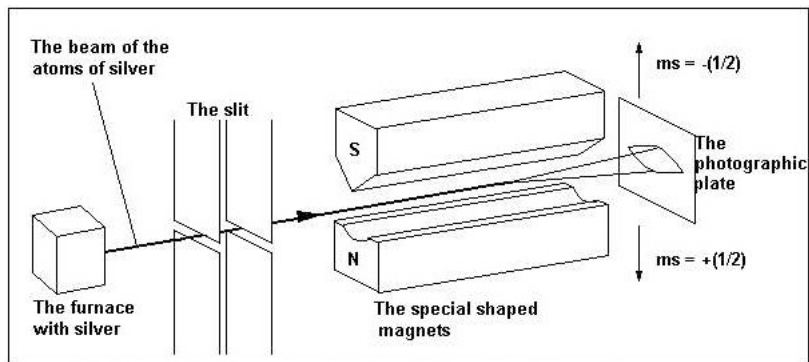
Force in a Field Gradient

$$\mathbf{F} = -\nabla E = \mu_z \frac{\partial B_z}{\partial z}$$



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Stern-Gerlach Experiment



The Stern-Gerlach experiment. On the photographic plate are two clear tracks.

Image from <http://library.thinkquest.org/19662/high/eng/exp-stern-gerlach.html?tqskip=1>

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Stern-Gerlach Experiment

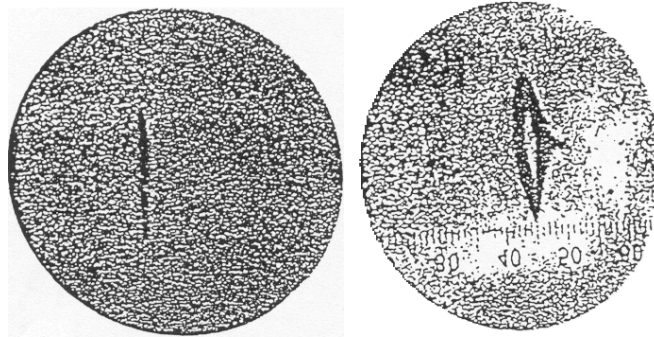


Image from <http://library.thinkquest.org/19662/high/eng/exp-stern-gerlach.html?tskip=1>

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Quantization of Magnetic Moment

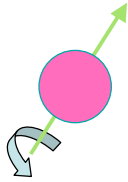
The key finding of the Stern-Gerlach experiment is that the magnetic moment is quantized. That is, it can only take on discrete values.

In the experiment, the finding was that

$$\mu_z = +\mu_0 \text{ OR } -\mu_0$$

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Magnetic Moment and Angular Momentum



A charged sphere spinning about its axis has angular momentum and a magnetic moment.

This is a classical analogy that is useful for understanding quantum spin, but remember that it is only an analogy!

Relation: $\boldsymbol{\mu} = \gamma \mathbf{S}$ where γ is the gyromagnetic ratio and \mathbf{S} is the spin angular momentum.

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Quantization of Angular Momentum

Because the magnetic moment is quantized, so is the angular momentum.

In particular, the z-component of the angular momentum is quantized as follows:

$$S_z = m_s \hbar$$

$$m_s \in \{-s, -(s-1), \dots, s\}$$

s is an integer or half integer

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Hydrogen Proton

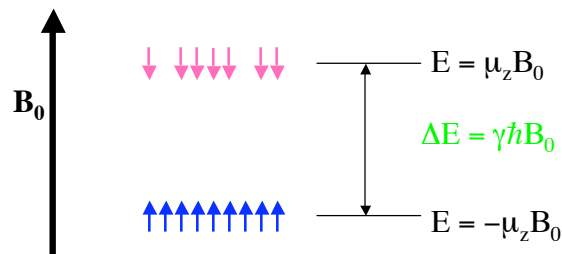
Spin 1/2

$$S_z = \begin{cases} +\hbar/2 \\ -\hbar/2 \end{cases}$$

$$\mu_z = \begin{cases} +\gamma\hbar/2 \\ -\gamma\hbar/2 \end{cases}$$

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Boltzmann Distribution



$$\frac{\text{Number Spins Up}}{\text{Number Spins Down}} = \exp(-\Delta E/kT)$$

Ratio = 0.999990 at 1.5T !!!

Corresponds to an excess of about 10 up spins per million

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Equilibrium Magnetization

$$\begin{aligned}\mathbf{M}_0 &= N\langle\mu_z\rangle = N\left(\frac{n_{up}(-\mu_z) + n_{down}(\mu_z)}{N}\right) \\ &= N\mu\frac{e^{\mu_z B/kT} - e^{-\mu_z B/kT}}{e^{\mu_z B/kT} + e^{-\mu_z B/kT}} \\ &\approx N\mu_z^2 B/(kT) \\ &= N\gamma^2\hbar^2 B/(4kT)\end{aligned}$$

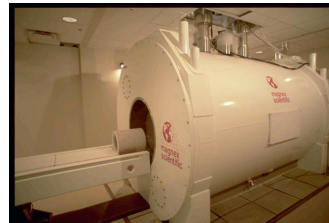
N = number of nuclear spins per unit volume
Magnetization is proportional to applied field.

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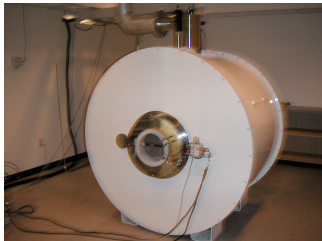
Bigger is better



3T Human imager at UCSD.



7T Human imager at U. Minn.



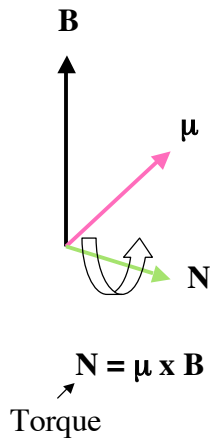
7T Rodent Imager at UCSD



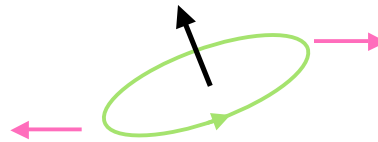
9.4T Human imager at UIC

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Torque



For a non-spinning magnetic moment, the torque will try to align the moment with magnetic field (e.g. compass needle)



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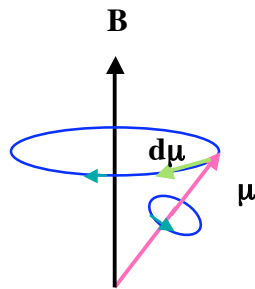
Precession

$$\begin{array}{l}
 \text{Torque} \\
 \downarrow \\
 N = \mu \times B \\
 \left. \begin{array}{l} N = \mu \times B \\ \frac{dS}{dt} = N \end{array} \right\} \\
 \uparrow \\
 \text{Change in} \\
 \text{Angular momentum}
 \end{array}
 \quad
 \begin{array}{l}
 \frac{dS}{dt} = \mu \times B \\
 \left. \begin{array}{l} \frac{dS}{dt} = \mu \times B \\ \mu = \gamma S \end{array} \right\} \\
 \uparrow \\
 \text{Relation between} \\
 \text{magnetic moment and} \\
 \text{angular momentum}
 \end{array}
 \quad
 \frac{d\mu}{dt} = \mu \times \gamma B$$

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Precession

$$\frac{d\boldsymbol{\mu}}{dt} = \boldsymbol{\mu} \times \gamma \mathbf{B}$$



Analogous to motion of a gyroscope

Precesses at an angular frequency of

$$\omega = \gamma B$$

This is known as the **Larmor** frequency.

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Larmor Frequency

$\omega = \gamma B$ Angular frequency in rad/sec

$f = \gamma B / (2\pi)$ Frequency in cycles/sec or Hertz,
Abbreviated Hz

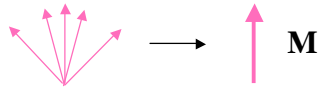
For a 1.5 T system, the Larmor frequency is 63.86 MHz which is 63.86 million cycles per second. For comparison, KPBS-FM transmits at 89.5 MHz.

Note that the earth's magnetic field is about 50 μT , so that a 1.5T system is about 30,000 times stronger.

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Magnetization Vector

$$\mathbf{M} = \frac{1}{V} \sum_{\substack{\text{protons} \\ \text{in } V}} \mu_i$$



Vector sum of the magnetic moments over a volume.

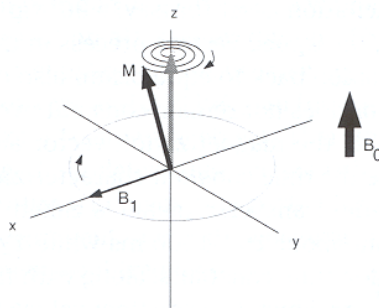
For a sample at equilibrium in a magnetic field, the transverse components of the moments cancel out, so that there is only a longitudinal component.

$$\frac{d\mathbf{M}}{dt} = \gamma \mathbf{M} \times \mathbf{B}$$

Equation of motion is the same form as for individual moments.

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RF Excitation



At equilibrium, net magnetization is parallel to the main magnetic field. How do we tip the magnetization away from equilibrium?

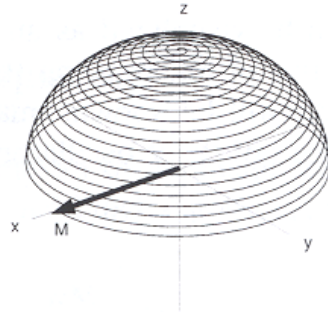
B_1 radiofrequency field tuned to Larmor frequency and applied in transverse (xy) plane induces nutation (at Larmor frequency) of magnetization vector as it tips away from the z -axis.

- lab frame of reference

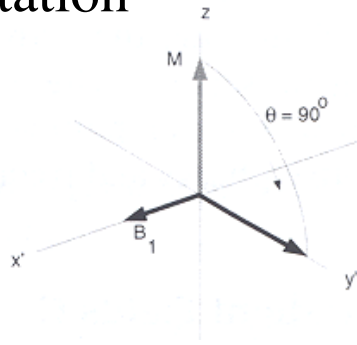
Image & caption: Nishimura, Fig. 3.2

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RF Excitation



a) Laboratory frame behavior of \mathbf{M}

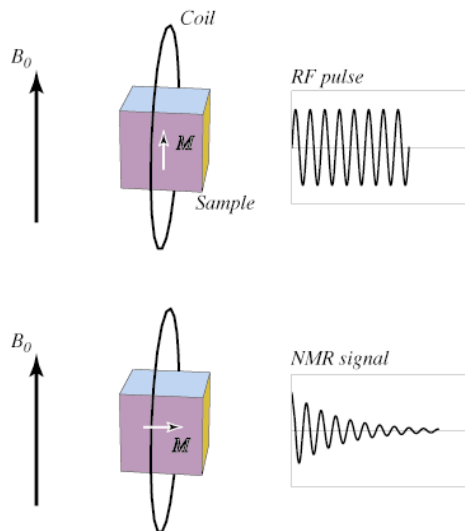


b) Rotating frame behavior of \mathbf{M}

B_1 induces rotation of magnetization towards the transverse plane. Strength and duration of B_1 can be set for a 90 degree rotation, leaving \mathbf{M} entirely in the xy plane. Images & caption: Nishimura, Fig. 3.3

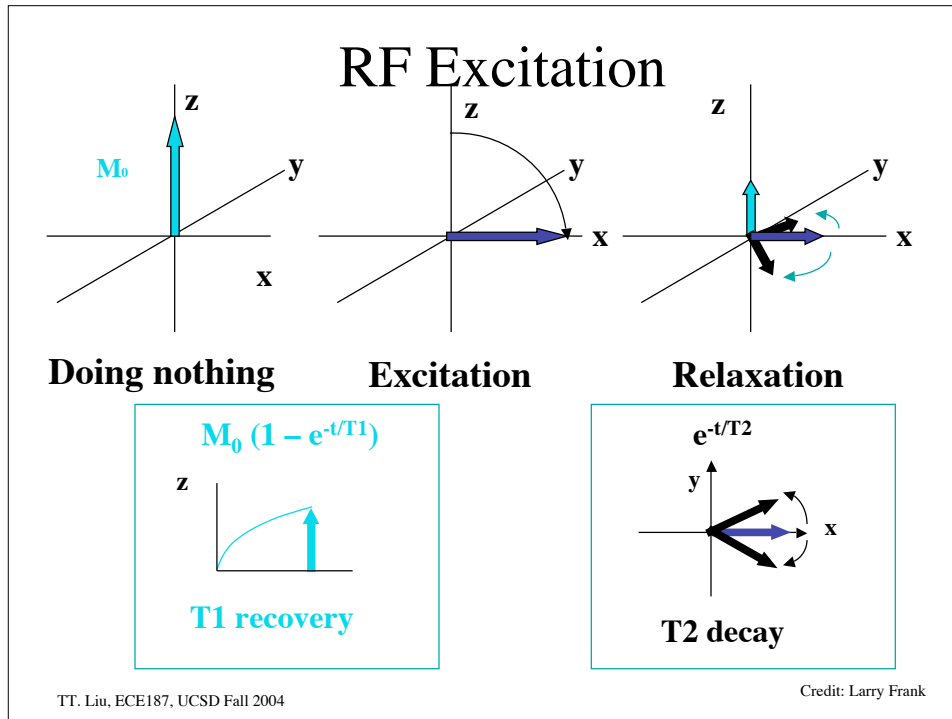
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RF Excitation



From Buxton 2002

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Relaxation

An excitation pulse rotates the magnetization vector away from its equilibrium state (purely longitudinal). The resulting vector has both longitudinal M_z and transverse M_{xy} components.

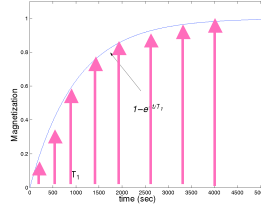
Due to thermal interactions, the magnetization will return to its equilibrium state with characteristic time constants.

- T_1 spin-lattice time constant, return to equilibrium of M_z
- T_2 spin-spin time constant, return to equilibrium of M_{xy}

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Longitudinal Relaxation

$$\frac{dM_z}{dt} = -\frac{M_z - M_0}{T_1}$$



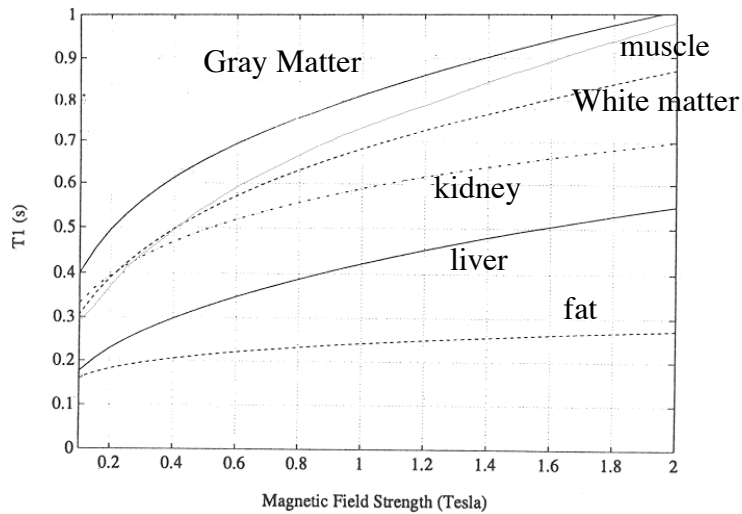
After a 90 degree pulse $M_z(t) = M_0(1 - e^{-t/T_1})$

Due to exchange of energy between nuclei and the lattice (thermal vibrations). Process continues until thermal equilibrium as determined by Boltzmann statistics is obtained.

The energy ΔE required for transitions between down to up spins, increases with field strength, so that T_1 increases with \mathbf{B} .

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T1 Values

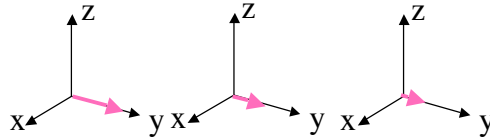


Image, caption: Nishimura, Fig. 4.2

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Transverse Relaxation

$$\frac{dM_{xy}}{dt} = -\frac{M_{xy}}{T_2}$$



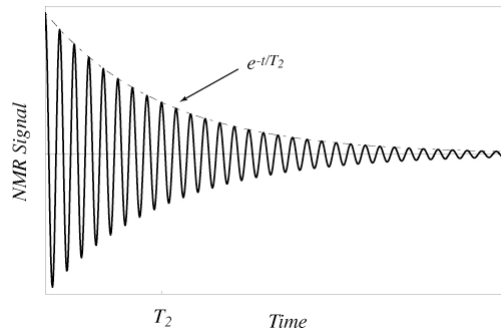
Each spin's local field is affected by the z-component of the field due to other spins. Thus, the Larmor frequency of each spin will be slightly different. This leads to a dephasing of the transverse magnetization, which is characterized by an exponential decay.

T_2 is largely independent of field. T_2 is short for low frequency fluctuations, such as those associated with slowly tumbling macromolecules.

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T2 Relaxation

Free Induction Decay (FID)

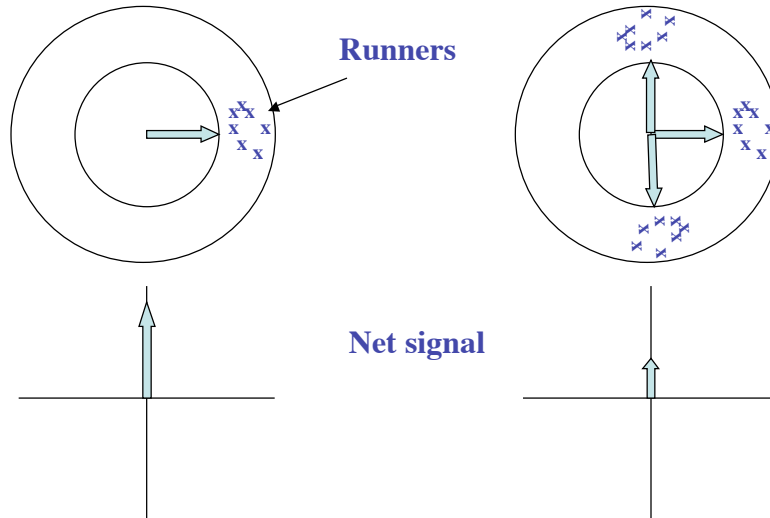


After a 90 degree
excitation

$$M_{xy}(t) = M_0 e^{-t/T_2}$$

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T2 Relaxation



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Credit: Larry Frank

T2 Values

Tissue	T_2 (ms)
gray matter	100
white matter	92
muscle	47
fat	85
kidney	58
liver	43
CSF	4000

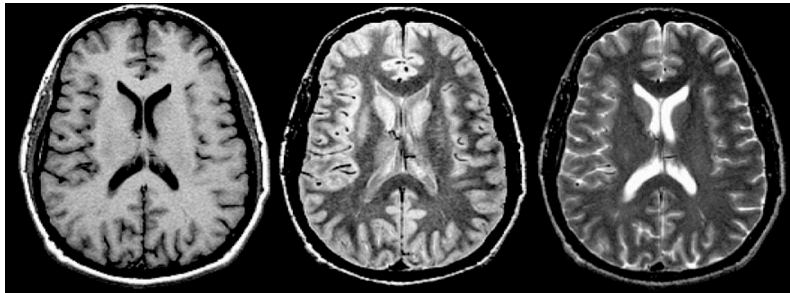
Solids exhibit very short T_2 relaxation times because there are many low frequency interactions between the immobile spins.

On the other hand, liquids show relatively long T_2 values, because the spins are highly mobile and net fields average out.

Table: adapted from Nishimura, Table 4.2

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Example



T₁-weighted

Density-weighted

T₂-weighted

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Bloch Equation

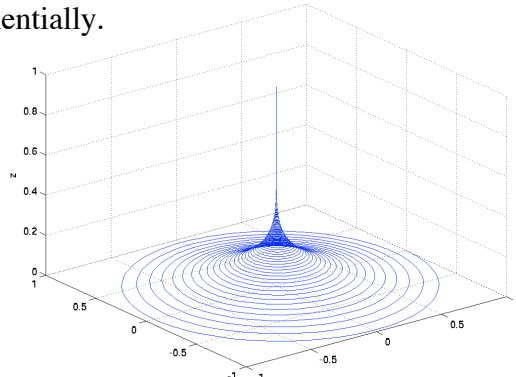
$$\frac{d\mathbf{M}}{dt} = \underbrace{\mathbf{M} \times \gamma \mathbf{B}}_{\text{Precession}} - \underbrace{\frac{M_x \mathbf{i} + M_y \mathbf{j}}{T_2}}_{\text{Transverse Relaxation}} - \underbrace{\frac{(M_z - M_0) \mathbf{k}}{T_1}}_{\text{Longitudinal Relaxation}}$$

$\mathbf{i}, \mathbf{j}, \mathbf{k}$ are unit vectors in the x,y,z directions.

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Summary

- 1) Longitudinal component recovers exponentially.
- 2) Transverse component precesses and decays exponentially.



Fact: Can show that $T_2 < T_1$ in order for $|M(t)| \leq M_0$
 Physically, the mechanisms that give rise to T_1 relaxation
 also contribute to transverse T_2 relaxation.

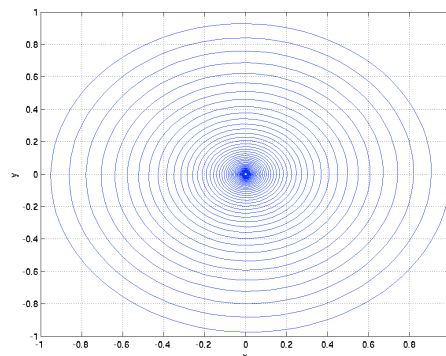
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Transverse Component

$$M \equiv M_x + jM_y$$

$$\begin{aligned} dM/dt &= d/dt(M_x + iM_y) \\ &= -j(\omega_0 + 1/T_2)M \end{aligned}$$

$$M(t) = M(0)e^{-j\omega_0 t} e^{-t/T_2}$$



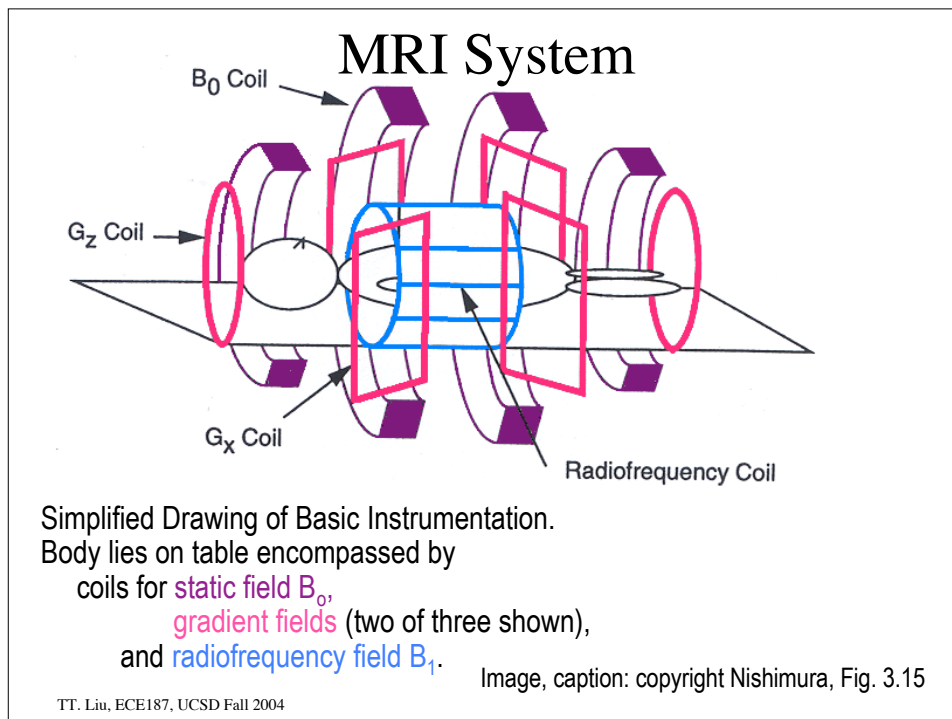
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Gradients

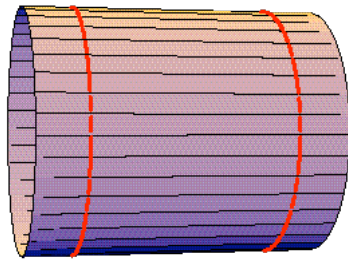
Spins precess at the Larmor frequency, which is proportional to the local magnetic field. In a constant magnetic field $B_z=B_0$, all the spins precess at the same frequency (ignoring chemical shift).

Gradient coils are used to add a spatial variation to B_z such that $B_z(x,y,z) = B_0 + \Delta B_z(x,y,z)$. Thus, spins at different physical locations will precess at different frequencies.

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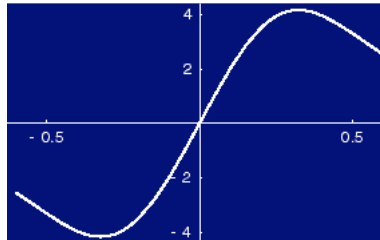


Z Gradient Coil



L

B(mT)



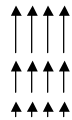
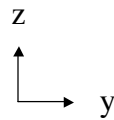
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Credit: Buxton 2002

Gradient Fields

$$B_z(x, y, z) = B_0 + \frac{\partial B_z}{\partial x} x + \frac{\partial B_z}{\partial y} y + \frac{\partial B_z}{\partial z} z$$

$$= B_0 + G_x x + G_y y + G_z z$$



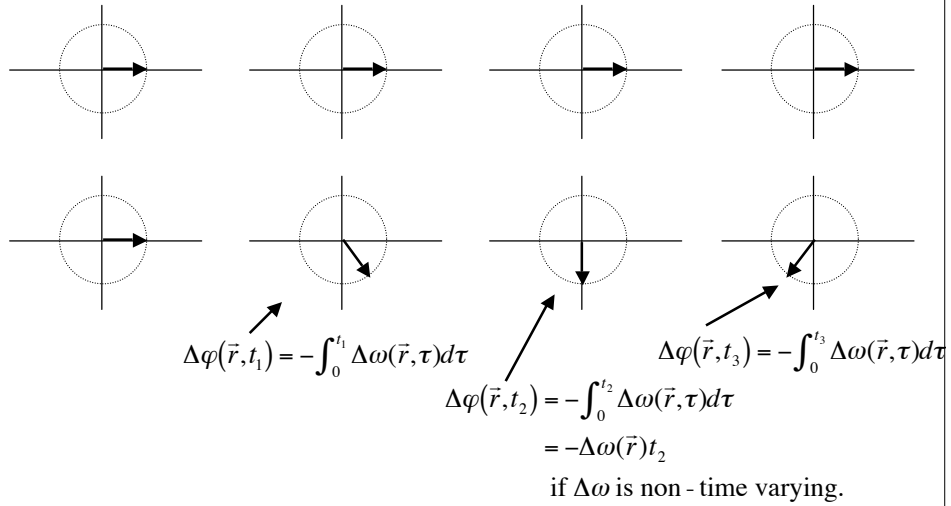
$$G_z = \frac{\partial B_z}{\partial z} > 0$$



$$G_y = \frac{\partial B_z}{\partial y} > 0$$

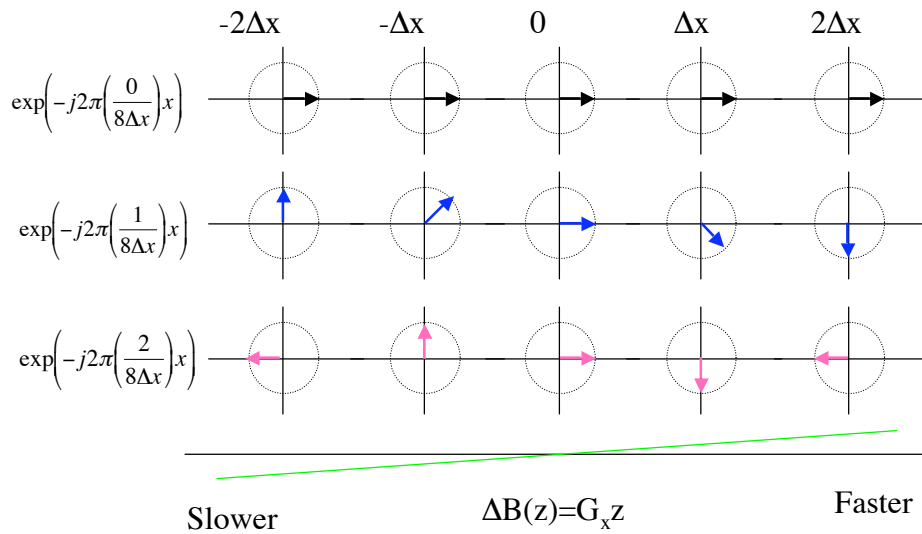
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Precession with Gradient Field



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Interpretation



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K-space

At each point in time, the received signal is the Fourier transform of the object

$$s(t) = M(k_x(t), k_y(t)) = F[m(x, y)]_{k_x(t), k_y(t)}$$

evaluated at the spatial frequencies:

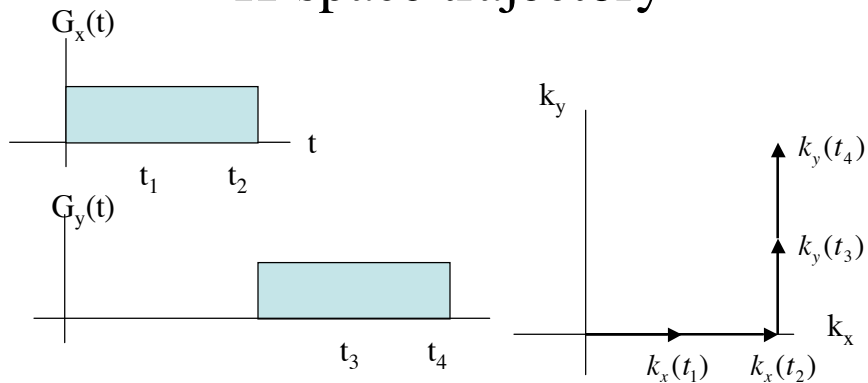
$$k_x(t) = \frac{\gamma}{2\pi} \int_0^t G_x(\tau) d\tau$$

$$k_y(t) = \frac{\gamma}{2\pi} \int_0^t G_y(\tau) d\tau$$

Thus, the gradients control our position in k-space. The design of an MRI pulse sequence requires us to efficiently cover enough of k-space to form our image.

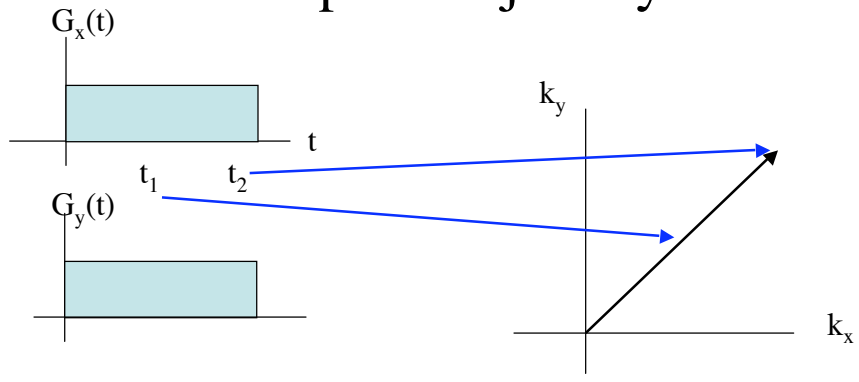
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K-space trajectory



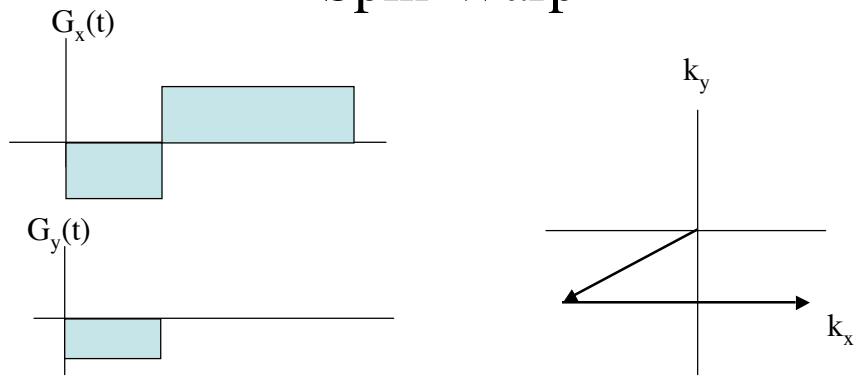
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K-space trajectory



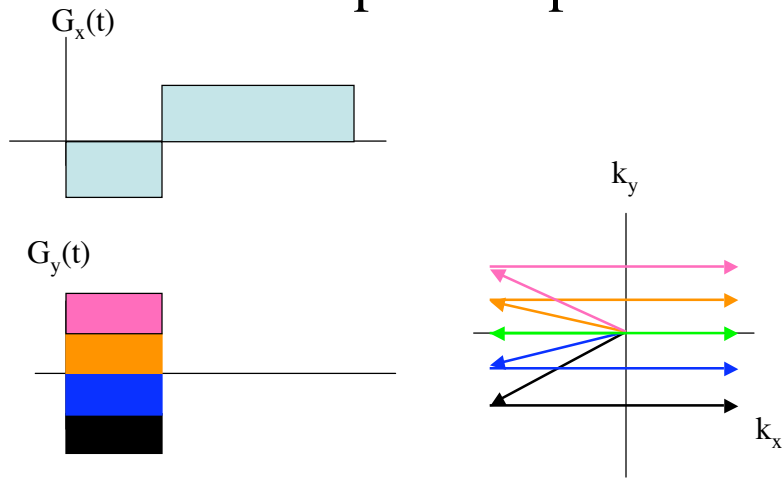
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Spin-Warp



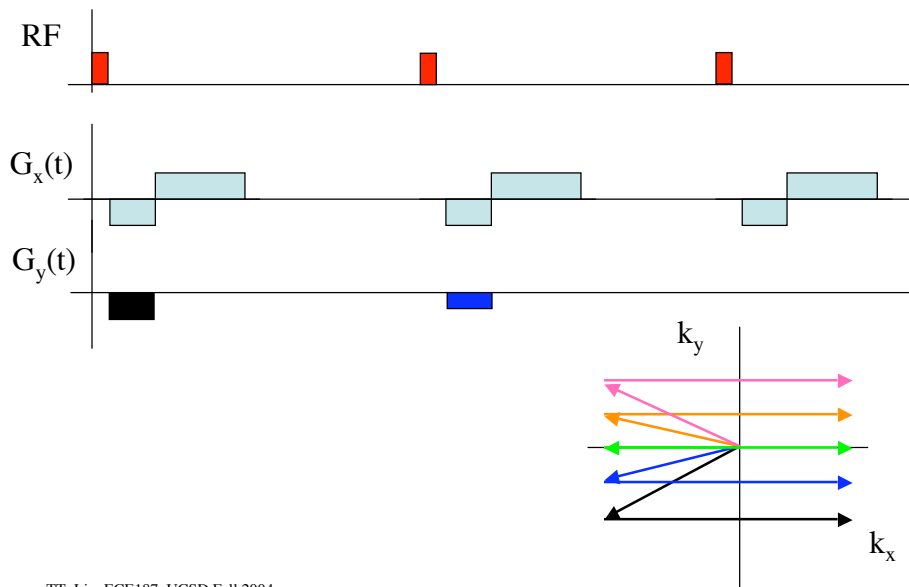
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Spin-Warp



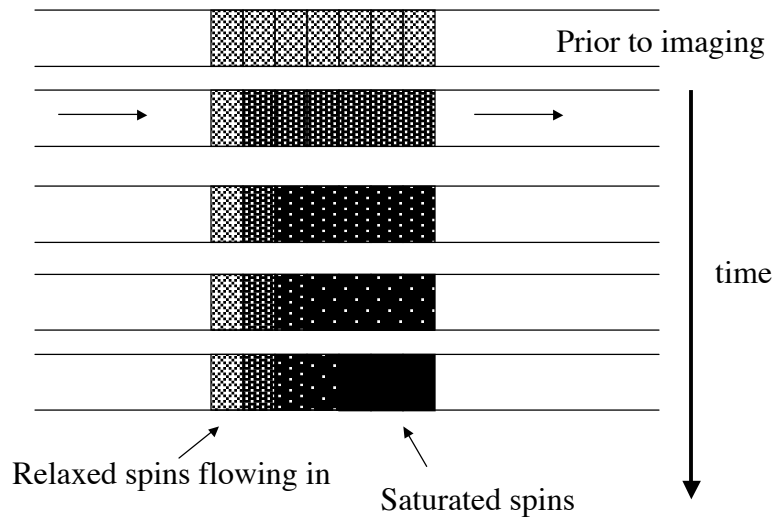
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Spin-Warp Pulse Sequence



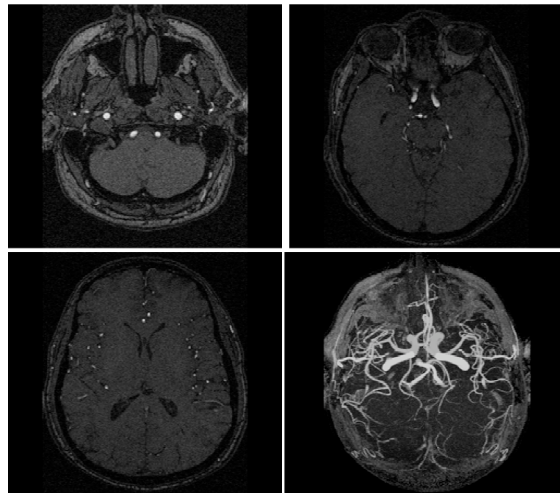
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Inflow Effect



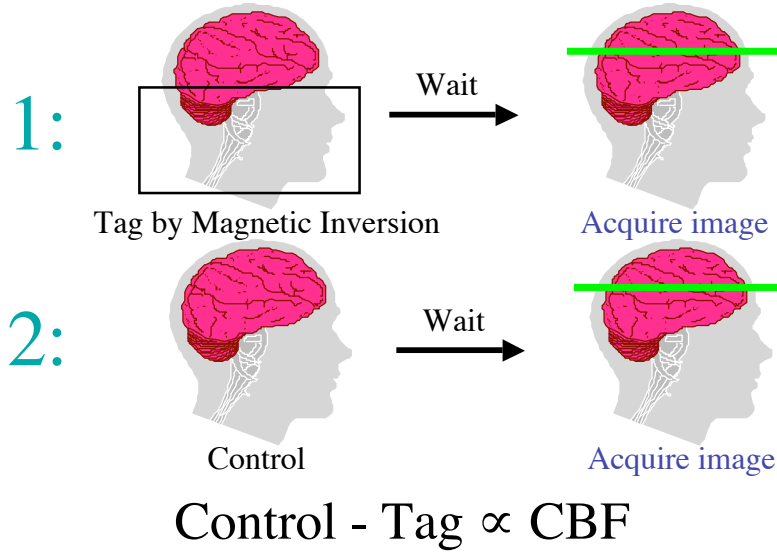
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Time of Flight Angiography



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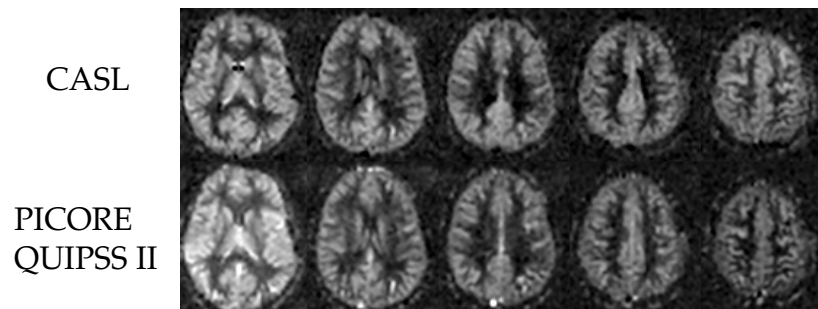
Arterial Spin Labeling (ASL)



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Credit: Wen-Ming Luh

Multislice CASL and PICORE



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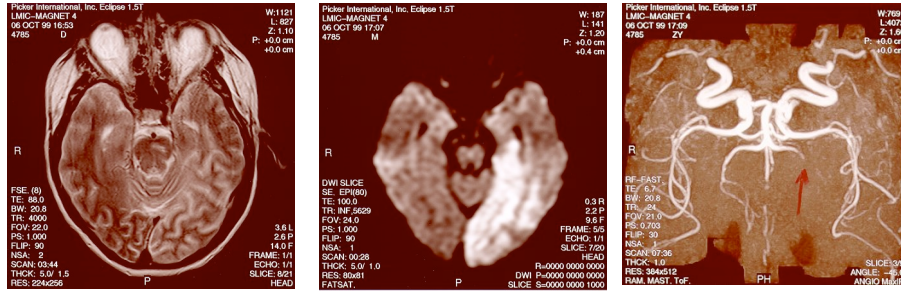
Credit: E. Wong

Diffusion Weighted Images

T2 weighted

Diffusion Weighted

Angiogram



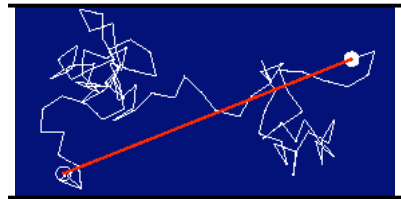
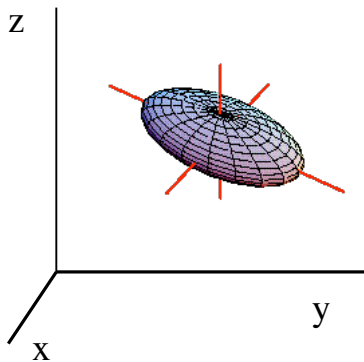
After a stroke, normal water movement is restricted in the region of damage. Diffusivity decreases, so the signal intensity increases.

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<http://lehighmri.com/cases/dwi/patient-b.html>

Restricted Diffusion

D depends on direction

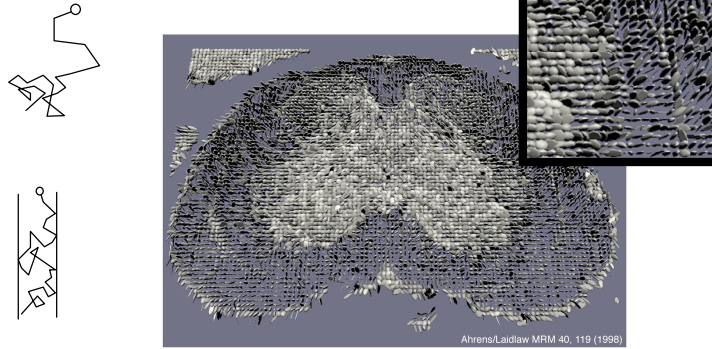


Diffusion tensor:
3 values of D
3 angles

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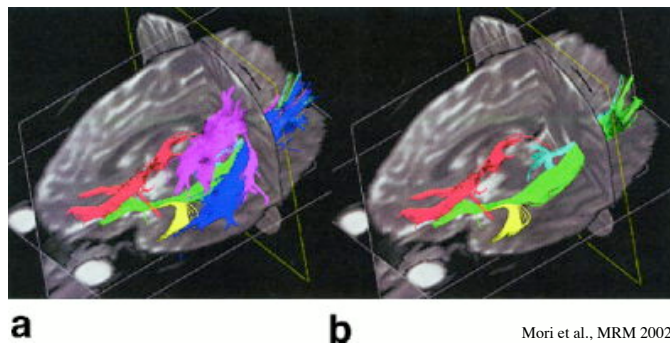
Credit: Larry Frank

Diffusion Imaging Example



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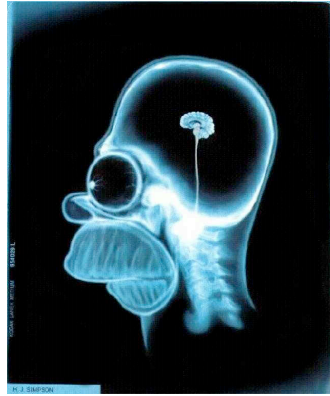
Fiber Tract Mapping



TT. Liu, ECE187, UCSD Fall 2004

fMRI

MRI studies brain anatomy.



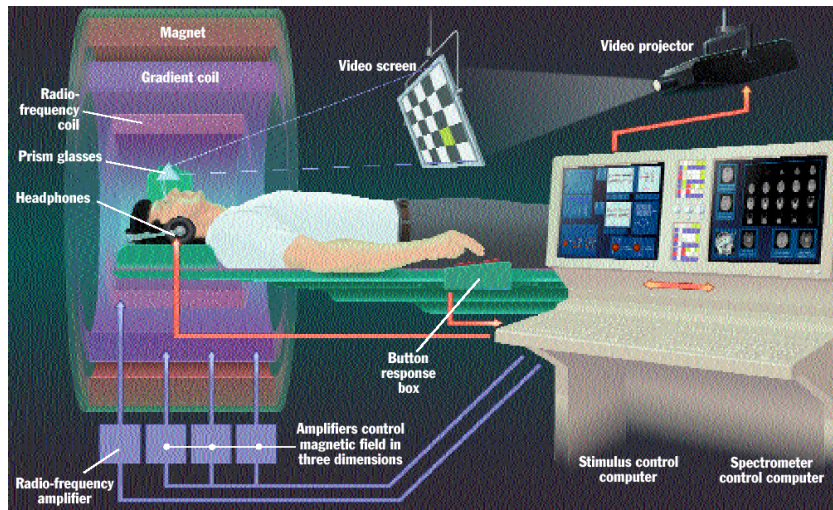
Functional MRI (fMRI) studies brain function.



TT. Liu, ECE187, UCSD Fall 2004

http://defiant.ssc.uwo.ca/Jody_web/fmri4dummies.htm

fMRI Setup



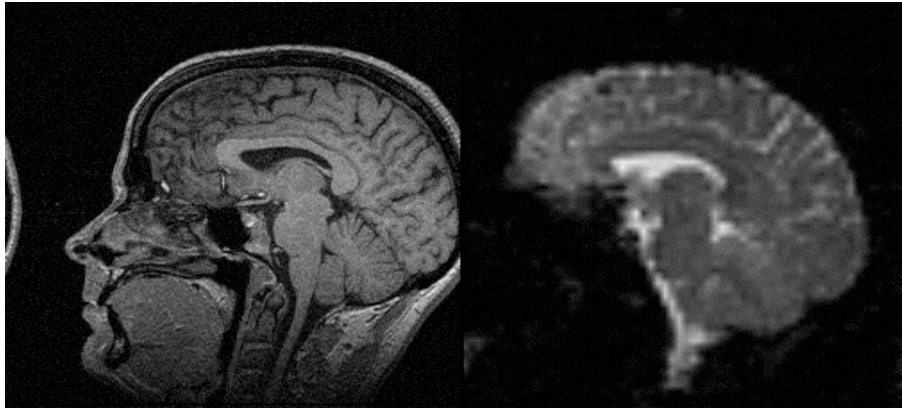
TT. Liu, ECE187, UCSD Fall 2004

http://defiant.ssc.uwo.ca/Jody_web/fmri4dummies.htm

fMRI Acquisition

High spatial resolution

High temporal resolution



MP-RAGE

Voxel volume: 1 mm³

Imaging time: 6 min

EPI

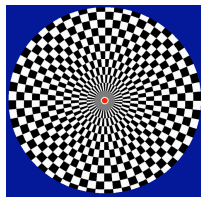
Voxel volume: 45 mm³

Imaging time: 60 msec

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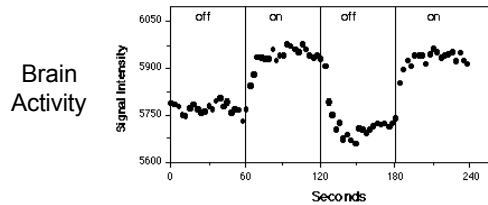
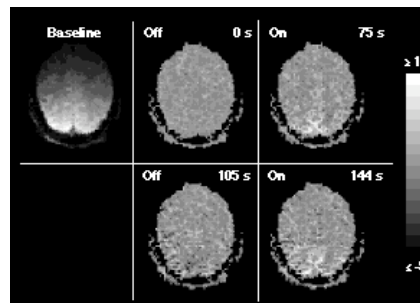
Buxton 2002

Visual Activation



Flickering Checkerboard

OFF (60 s) - ON (60 s) - OFF (60 s) - ON (60 s) - OFF (60 s)



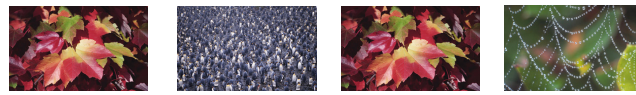
Time \Rightarrow

Source: Kwong et al., 1992

TT. Liu, ECE187, UCSD Fall 2004

http://defiant.ssc.uwo.ca/Jody_web/fmri4dummies.htm

Memory Encoding



10 Familiar scenes

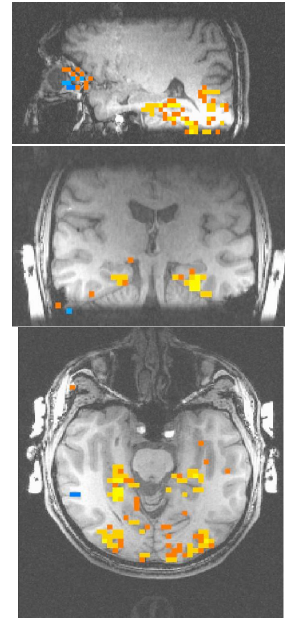
10 Novel Scenes

10 Familiar scenes

10 Novel Scenes

25 seconds

The same 4 familiar scenes are used throughout the experiment



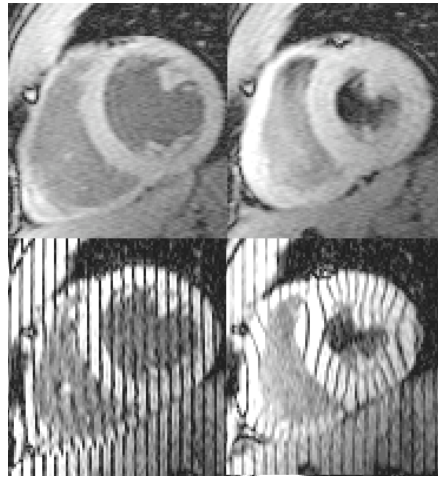
TT. Liu, ECE187, UCSD Fall 2004

Cardiac Imaging



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Cardiac Tagging



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